Two dimensional quantum memories Commuting projector codes

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Outline

- 2D Commuting Projector Codes
- Polographic Disentangling Lemma
- 3 Holographic Minimum Distance
 - 4 Capacity-Stability Tradeoff
- 5 String-Like Logical Operators
- Open Questions

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1 2D Commuting Projector Codes

- 2 Holographic Disentangling Lemma
- 3 Holographic Minimum Distance
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- 6 Open Questions

A is a 2D lattice.

- Each vertex occupied by d-level quantum particle.
- Hamiltonian $H = -\sum_{X \subset \Lambda} P_X$ with
 - $P_X = 0$ if radius $(X) \ge w$.
 - $[P_X, P_Y] = 0.$
 - *P_X* are projectors (optional).
- Code $C = \{\psi : P_X | \psi \rangle = |\psi \rangle \}$
 - = ground space of *H*
 - = image of code projector $\Pi = \prod_X P_X$
- With proper coarse graining, we can assume that
 - A is a regular square lattice.
 - Each P_X acts on 2 × 2 cell.

- Λ is a 2D lattice.
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2D Commuting Projector Codes Well known examples

Kitaev's toric code

- Bombin's topological color codes
- Levin & Wen's string-net models
- Turaev-Viro models
- Kitaev's quantum double models
- Most known models with topological quantum order

Remark

The first two example are simple because they are stabilizer codes. Most things I will say are trivial to prove in this case.

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Correctable region

A region $M \subset \Lambda$ is *correctable* if there exists a recovery operation \mathcal{R} such that $\mathcal{R}(\operatorname{Tr}_M \rho) = \rho$ for all code states ρ .

Minimum distance

The minimum distance d is the size of the smallest non-correctable region.

Logical operator

Operator L such that $L|\psi\rangle$ is a code state for any code state $|\psi\rangle$.

Rate (capacity)

The rate of a code is $R = \frac{k}{n}$ where $k = \log \dim(\mathcal{C})$ and $n = |\Lambda|$ in the number of particles.

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Holographic disentangling lemma

Let $M \subset \Lambda$ be a correctable region and suppose that its boundary ∂M is also correctable. Then, there exists a unitary operator $U_{\partial M}$ acting only on the boundary of M such that, for any code state $|\psi\rangle$,

$$|\mathcal{J}_{\partial M}|\psi
angle = |\phi_M
angle \otimes |\psi'_{\overline{M}}
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for some *fixed* state $|\phi_M\rangle$ on *M*.

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For a trivial code k = 0, every region is correctable, so we recover the area law $S(M) \le |\partial M|$ for commuting Hamiltonians of Wolf, Verstraete, Hastings, and Cirac.

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• Let *M* be correctable.

- Assume ∂M is correctable.
- Let $M = A \cup B$, $\overline{M} = C \cup D$, and $\partial M = B \cup C$.
- Write $\Pi = P_{AB}P_{B\overline{M}}$ with $[P_{AB}, P_{B\overline{M}}] = 0$.



- $\mathcal{H}_B = \bigoplus_J \mathcal{H}_{B_I^{\prime}} \otimes \mathcal{H}_{B_B^{\prime}}$ and $\Pi = \bigoplus_J P_{AB_I^{\prime}} \otimes P_{B_B^{\prime}\overline{M}}$
- This last sum over J contains only one non-zero factor since B ⊂ M is correctable.
- We can divide *B* into two subsystems *B*¹ and *B*² such that $\Pi = V_B P_{AB^1} \otimes P_{B^2 \overline{M}} V_B^{\dagger}$. (*)

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- Combining (*) with (**), $\Pi' = V_B^{\dagger} V_C^{\dagger} \Pi V_B V_C = P_{AB^1} P_{B^2 C^1} P_{C^2 D}$
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- Then $U_{\partial M} = V_{B^2C^1} V_B^{\dagger} V_C^{\dagger}$ disentangles region *M* as claimed.

- Let *M* be correctable.
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- Combining (*) with (**), $\Pi' = V_B^{\dagger} V_C^{\dagger} \Pi V_B V_C = P_{AB^1} P_{B^2 C^1} P_{C^2 D}$
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Holographic Minimum Distance

Statement of the result

Holographic minimum distance

Region $M \subset \Lambda$ is correctable if its boundary is smaller than the minimum distance $|\partial M| \leq cd$.

- Bulky errors are not problematic: it's the skinny ones we need to worry about.
- This hints at our next result: string-like logical operators.

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• If $|\partial M| \leq d$, then ∂M is also correctable.



- Thus, we can reconstruct any code state ρ from $\rho_{AD} = \text{Tr}_{\partial M} \rho$.
- But from the Holographic disentangling lemma, $\rho_{AD} = \eta_A \otimes \rho_D$ with η_A independent of the encoded state ρ .
- Thus, we can reconstruct ρ from $\rho_D = \text{Tr}_{M \cup \partial M} \rho$, so $M \cup \partial M$ is correctable.
- We can continue to grow *M* this way until $|\partial M| \ge d$.

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Capacity-Stability Tradeoff

$$k \leq c \frac{n}{d^2}$$

- Singleton's bound: $k \le n 2(d 1)$.
- Hamming bound: $k \le n \left| 1 \frac{d}{2n} \log 3 H(\frac{d}{2n}) \right|$.
- Kitaev's codes (with punctures) saturate this bound, so it is tight.
 - Holes of linear size *l* separated by distance *l*.
 - Minimum distance d oc l.
 - Number of logical qubits $l < \infty$ number of holes $\infty n/\ell^2 \propto n/d^2$.
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Information-theoretic condition for error correction

M is correctable iff $S(M\overline{M}) = S(\overline{M}) - S(M)$ for any code state ρ .

Obvious for pure states.

- Let $\rho_{M\overline{M}}$ be a code state and $\rho_{M\overline{M}R}$ its purification.
- By assumption, there exists *R* on Λ such that *R*(Tr_Mρ_{MMR}) = ρ_{MMR} and *R*(Tr_Mρ_{MM} ⊗ ρ_R) = ρ_{MM} ⊗ ρ_R
- Since relative entropy can only decrease under the action of a CPTP map, $S(\rho_{M\overline{M}R} \| \rho_{M\overline{M}} \otimes \rho_R) = S(\text{Tr}_M \rho_{M\overline{M}R} \| \text{Tr}_M \rho_{M\overline{M}} \otimes \rho_R)$
- Using $S(\rho_{AB} || \rho_A \otimes \rho_B) = S(A) + S(B) S(AB)$ and the fact that $\rho_{M\overline{M}B}$ is pure, we get the desired result.
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Union of correctable regions

- Trivial for syndrome-based error correction (e.g. stabilizer codes).
 We will prove the Knill-Laflamme condition DOW & OV D < D
- The holographic disentangling lemma applied to M_1 implies that $\Pi = V_B V_C |\eta_{AB1}\rangle \langle \eta_{AB1} | \otimes |\nu_{B^2C1}\rangle \langle \nu_{B1C1} | \otimes P_{C^2D} V_B^{\dagger} V_C^{\dagger}$.
- So $\Pi O_{M1} \otimes O_{M_2} \Pi = f(O_{M_1}) \Pi O_{M_2} \Pi \propto \Pi$ where $f(O_{M_1}) = \langle \eta_{AB^1} | \langle \nu_{B^2 C^1} | V_B^{\dagger} O_{M_1} V^B | \eta_{AB^1} \rangle | \nu_{B^2 C^1} \rangle$.

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Proof of Capacity-Stability Tradeoff

- Squares have perimeter $\approx d$.
- Region *A* (union of blue squares) is correctable.
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 Let's apply the information theoretic conditions to maximally mixed code state ρ = Π/Tr(Π) in two different ways:

• Using $S(BC) \leq S(B) + S(C)$ and $S(AC) \leq S(A) + S(C)$ $2S(ABC) \leq 2S(C)$

S(ABC) = k. S(C) ≤ |C| ∝ number of circles ∝ n/d².

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Proof of Capacity-Stability tradeoff

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- 2D Commuting Projector Codes
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- 3 Holographic Minimum Distance
- 4 Capacity-Stability Tradeoff
- 5 String-Like Logical Operators

6 Open Questions

String-like logical operators

- Well known for Kitaev's toric code.
- Intuitive for known models that support anyons:
 - The ground state can be changed by dragging an anyon around a topologically non-trivial loop.
 - This process is realized on a string, and generated a logical operation.
- Relation to thermal instability?

String-like logical operators

There exists a non-trivial logical operator supported on a string-like region.

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- Let *M* be a string-like region.
- Suppose *M* is correctable.
- Consider its boundary $\partial M = \partial M_L \cup \partial M_R$.
- If either ∂M_L or ∂M_R are not correctable, we are done.
- Otherwise ∂M = ∂M_L ∪ ∂M_R is correctable, and therefore M ∪ ∂M is correctable.
- Continue until we arrive at A is correctable, which is impossible.

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• Let *M* be a non-correctable string-like region.

- There exists O_M such that $\Pi O_M \Pi \propto \Pi$.
- Let $\Pi_M = \prod_{X \cap M \neq \emptyset} P_X$
- Then $X = \prod_M O_M \prod_M$ is a non-trivial logical operator supported on $M \cup \partial M$.



 Any function of X, e.g. exp(-iXθ), is also a logical operator with the same support.

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• All our results extend to *D*-dimensional lattices, e.g. $k \le cn/d^{\frac{2}{D-1}}$

• How about infinite dimensions (LDPC codes)? (Delfosse & Zémor)

• String-like logical operators \Rightarrow constant energy barrier.

- This is not directly related to thermal instability.
- 2D Ising model has an energy barrier ∝ √n, but an energy ∝ n at finite temperature.
- What matters is entropy (for a given energy, many more configurations many with small error droplets than with a large one).
- Can we characterize all string-like logical operators'
- Relation between commuting projector codes and anyon models.
- Extend to frustration-free Hamiltonians (and therefore to all gapped Hamiltonians, i.e. Hastings).
 - Use proof techniques to show that gapped Hamiltonian \in QCMA.
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