



Fun with QMA and non-unitary gates



Dan Browne

joint work with: **Nairi Usher**



QMA

- ❖ Kitaev (1999) - **QMA**: The quantum analogue of **NP**

QMA for beginners



*talk
contains
Interstellar
spoilers - sorry!*

A FILM BY CHRISTOPHER NOLAN
INTERSTELLAR

NP



Arthur

Poly-time classical
computer



Merlin

Unbounded
computational
power



Arthur

NP



Merlin

Decision problem:

*“Is there an answer to
the **problem of gravity?**”*

Arthur can **verify**
validity of proof
in **poly-time**



*“**Yes**”*
*“And here’s the **proof**”*



Arthur

NP



Merlin

Decision problem:

*“Is there an answer to
the **problem of gravity?**”*

Arthur can **verify**
validity of proof
in **poly-time**

“No”

*“And there is **no**
fake proof I could
send to you to **trick**
you into thinking the
answer is yes”*



Arthur

Poly-time classical
computer

Decision problem

*Arthur can **verify**
validity of proof
in **poly-time***

NP



Merlin

Unbounded
computational
power

proof



“Yes = 1”

“And here’s the **proof**”

“No = 0”

***“And there is **no**
fake proof ...”***



Arthur

Poly-time classical computer

Decision problem

Arthur can **verify** validity of proof in **poly-time**

~~NP~~ MA



Merlin

Unbounded computational power

proof



with “**high enough**” probability

“**Yes = 1**”

“And here’s the **proof**”

“**No = 0**”

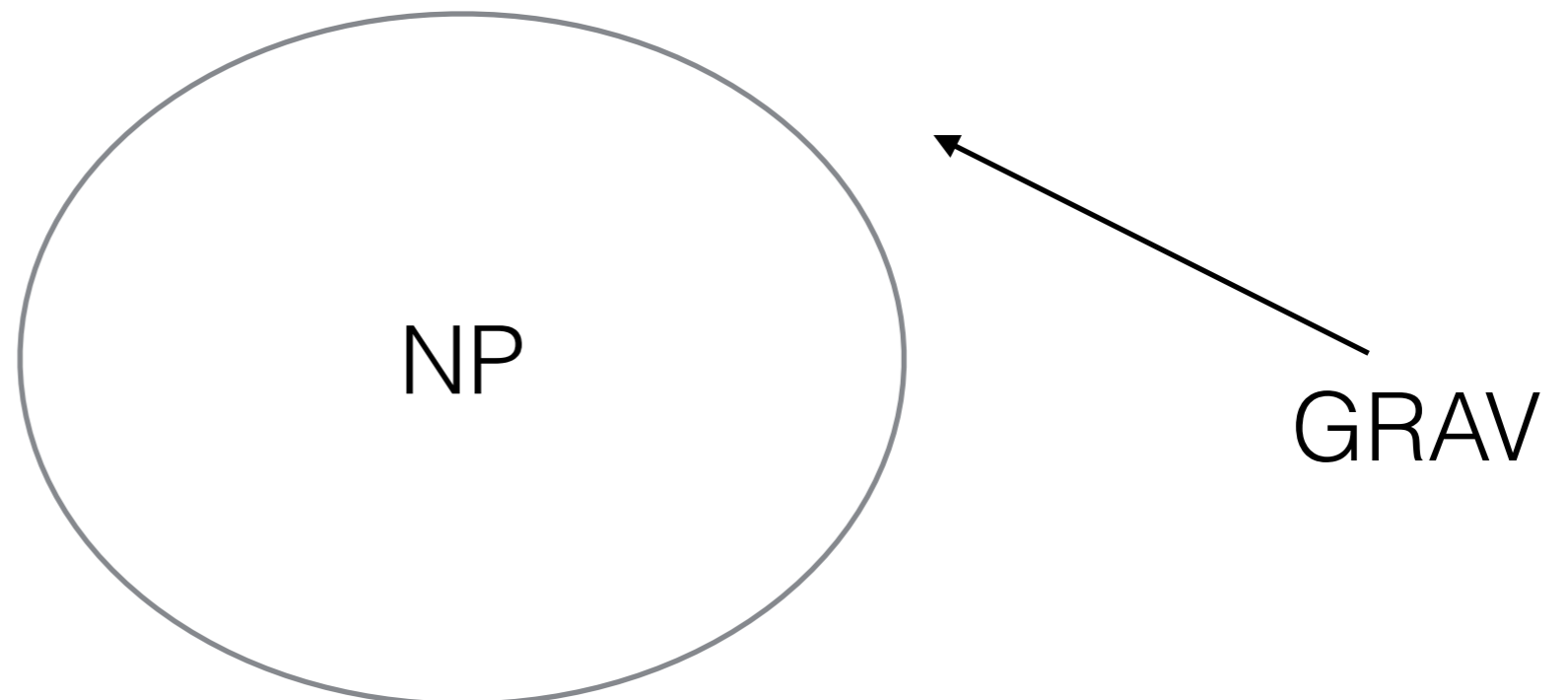
“And there is **no fake proof ...**”

*“Is there an answer to the **problem of gravity**?”*



*SPOILER: “I lied.
It is not in NP.”*

*We need **quantum data!***



QMA

- ❖ Kitaev (1999) - **QMA**: The quantum analogue of **NP**

QMA



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Poly-time **quantum**
computer

Decision problem

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Merlin

Unbounded
computational
power

quantum proof
←
 $|\psi\rangle$

with “**high
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probability

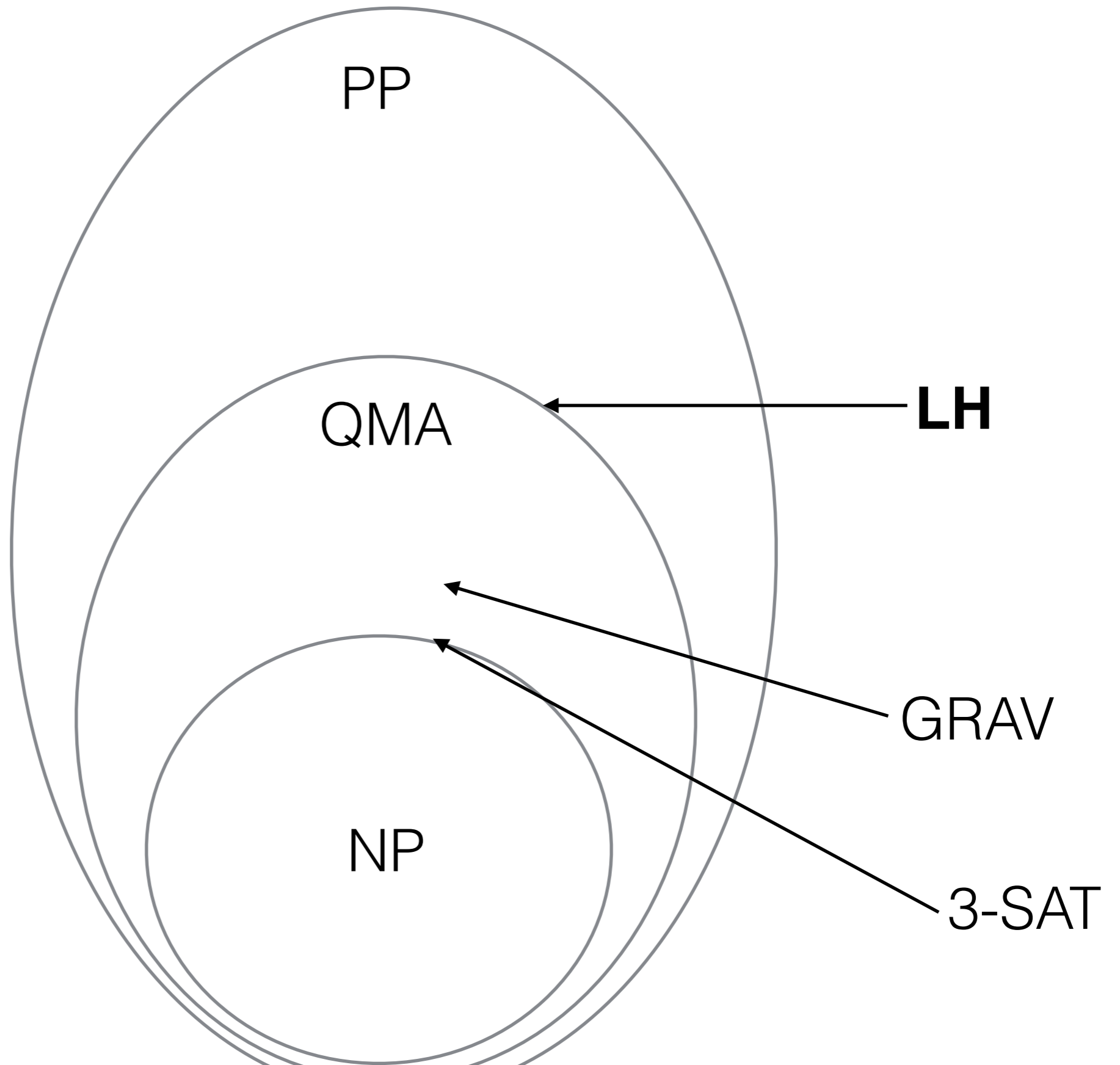
“**Yes = 1**”
“And here’s the **proof**”

“**No = 0**”
“And there is **no
fake proof ...**”

Local Hamiltonian Problem

Kitaev
(1999):

LH: k-local
Hamiltonian
Problem is
QMA-
complete



Local Hamiltonian Problem

Given: A k -local n -qubit Hamiltonian

$$H = \sum_j H_j$$

$$0 \leq H_j \leq 1$$

$$H_j : k\text{-local}$$

with a promise that the **ground state energy** of H is:

$$E_0 > b$$

or

$$0 < E_0 < a$$

$$E = b \text{ -----}$$

$$E = a \text{ -----}$$

$$E = 0 \text{ -----}$$

“promise gap” $> 1/\text{poly } n$

Problem: Determine whether $E_0 > b$ or $E_0 < a$?

LH is QMA-complete

LH is **in** QMA

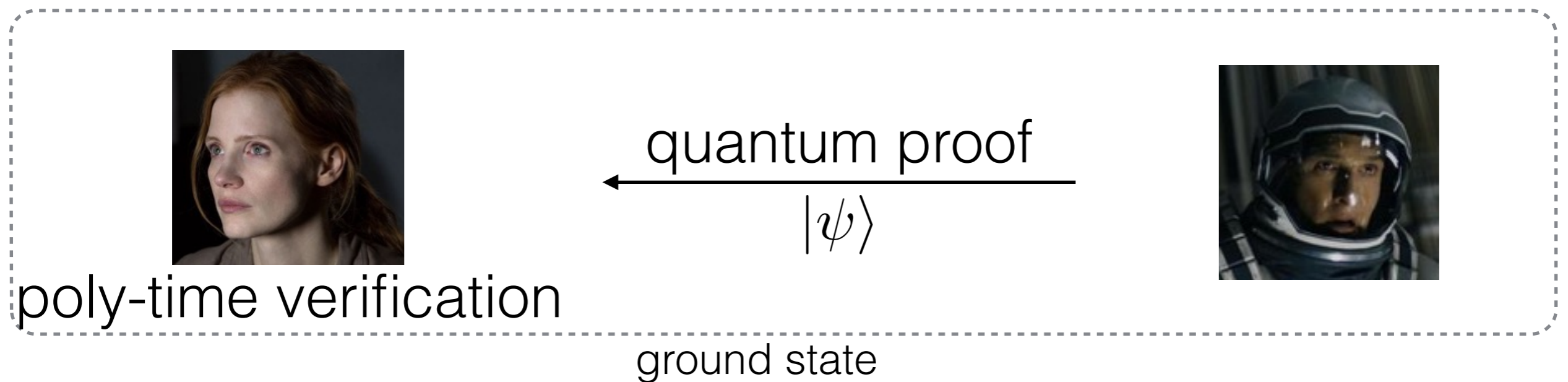


ground state of H
←
 $|\psi\rangle$

Can efficiently estimate
the energy on a
quantum computer.

LH is **hard** for QMA

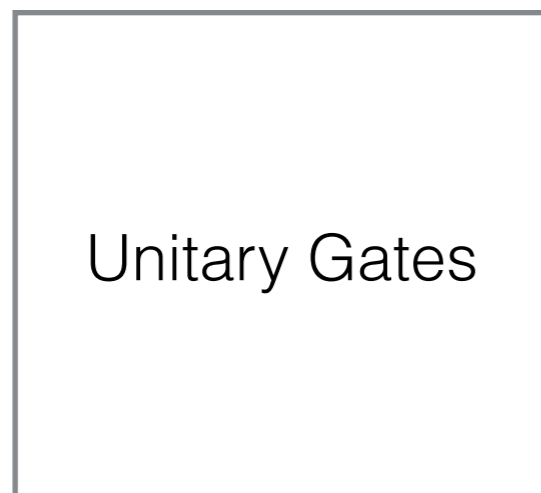
Encode proof verification for any QMA problem into a ground state of a LH



Yes instances: **Low** energy
No instances: **Higher** energy

Encoding a computation in a ground state

Feynman's **History state** idea



History State

$$\begin{aligned} |\psi\rangle = & |\psi_{\text{in}}\rangle |0\rangle \\ & + U_1 |\psi_{\text{in}}\rangle |1\rangle \\ & + U_2 U_1 |\psi_{\text{in}}\rangle |2\rangle \\ & + U_3 U_2 U_1 |\psi_{\text{in}}\rangle |3\rangle \\ & + \dots \end{aligned}$$

Encoding a computation in a ground state

Propagator Hamiltonian:

$$H_{\text{prop}} = \sum_j H_j \quad H_j = -\frac{1}{2} (U_j \otimes |j\rangle\langle j-1| + \text{h.c.}) + \frac{1}{2} \mathbb{1} \otimes (|j\rangle\langle j| + |j+1\rangle\langle j+1|)$$

**Ground space is
space of history states:**

$$\begin{aligned} & |\psi_{\text{in}}\rangle |0\rangle \\ & + U_1 |\psi_{\text{in}}\rangle |1\rangle \\ & + U_2 U_1 |\psi_{\text{in}}\rangle |2\rangle \\ & + U_3 U_2 U_1 |\psi_{\text{in}}\rangle |3\rangle \\ & + \dots \end{aligned}$$

Encoding a computation in a ground state

This works, because the **ground space** of

$$H_j = -\frac{1}{2} (U_j \otimes |j\rangle\langle j-1| + \text{h.c.}) + \frac{1}{2} \mathbb{1} \otimes (|j\rangle\langle j| + |j+1\rangle\langle j+1|)$$

is:

$$|\psi\rangle|j-1\rangle + U_j|\psi\rangle|j\rangle$$

Encoding a computation in a ground state

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$$H_j = -\frac{1}{2} (U_j \otimes |j\rangle\langle j-1| + \text{h.c.}) + \frac{1}{2} \mathbb{1} \otimes (|j\rangle\langle j| + |j+1\rangle\langle j+1|)$$

is:

$$\begin{aligned} & H_j(|\psi\rangle|j-1\rangle + U_j|\psi\rangle|j\rangle) \\ &= (1/2)(-U_j|\psi\rangle|j\rangle + U_j|\psi\rangle|j\rangle) = 0 \end{aligned}$$

QMA



Arthur

Poly-time **quantum**
computer

Decision problem

*Arthur can **verify**
validity of proof
in **poly-time***



Merlin

Unbounded
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power

quantum proof
←
 $|\psi\rangle$

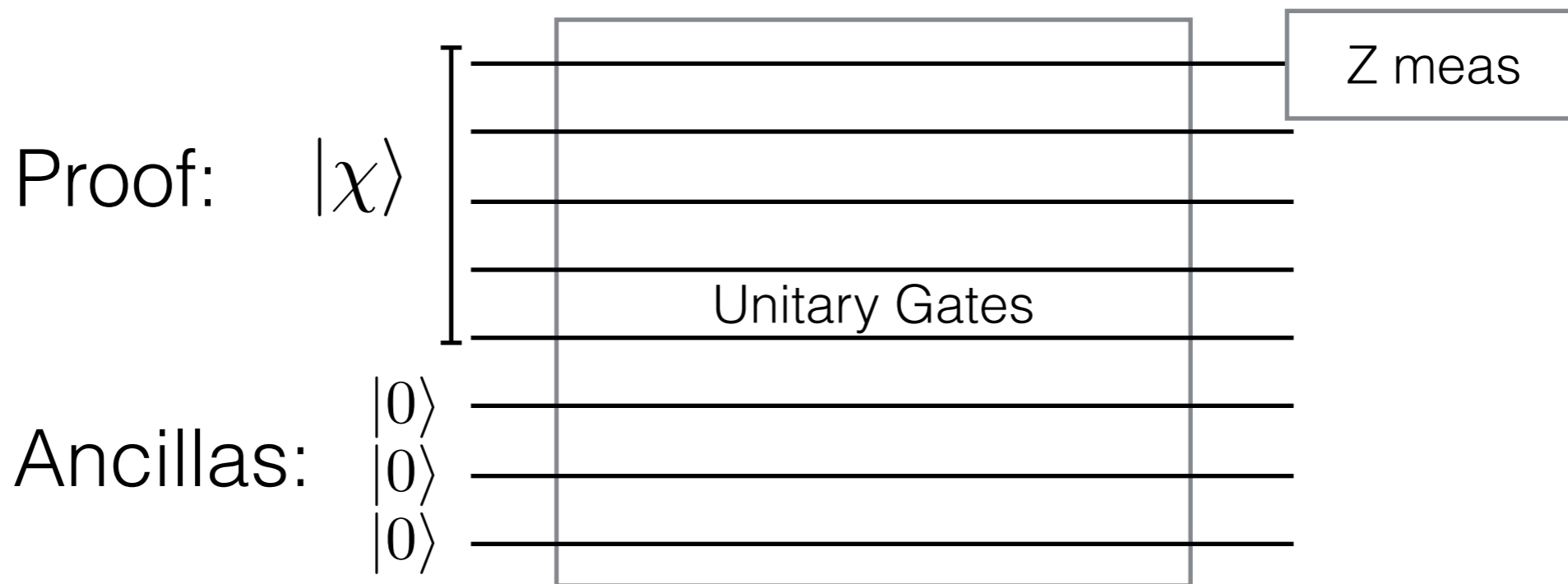
with “**high
enough**”
probability

“**Yes = 1**”
“And here’s the **proof**”

“**No = 0**”
“And there is **no
fake proof ...**”

Encoding a computation in a ground state

The verification circuit we wish to encode:

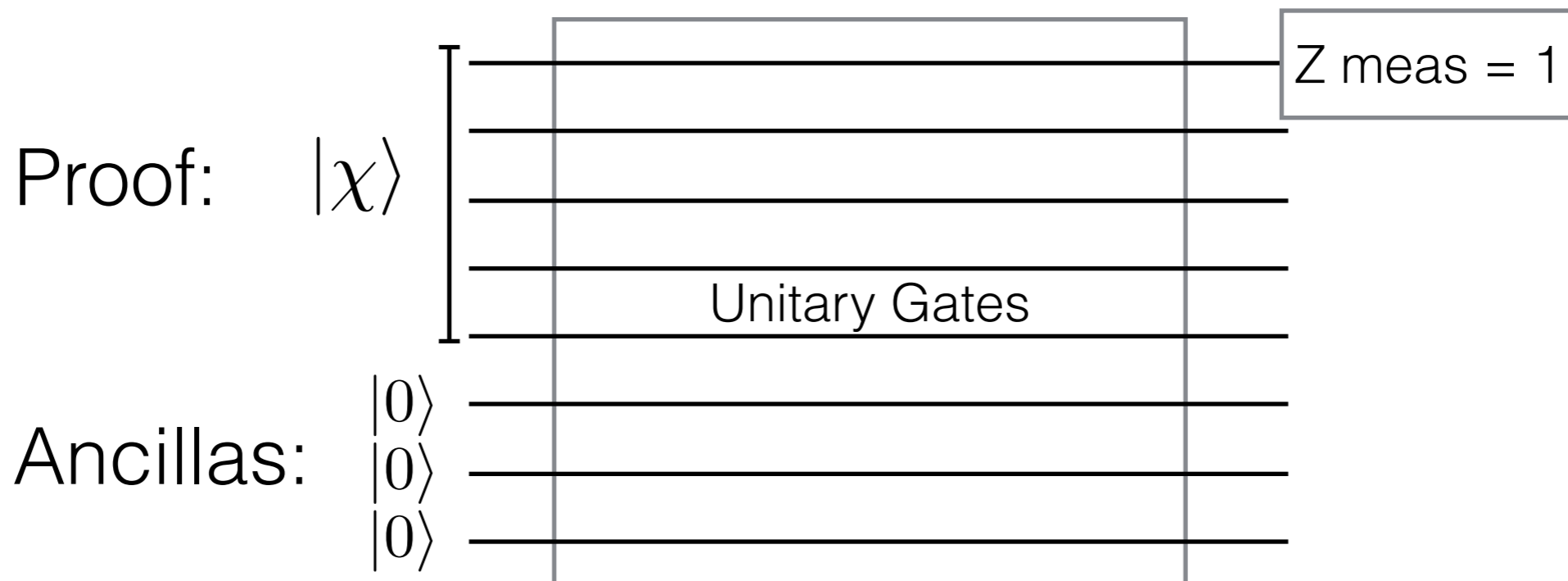


L = No. of unitary gates in circuit

Encoding a computation in a ground state

Maps into the **ground space** of

$$H = H_{\text{in}} + H_{\text{prop}} + H_{\text{out}}$$



$$H_{\text{in}} = \sum_{j \in \text{ancilla}} |1\rangle_j \langle 1| \otimes |0\rangle_{\text{clock}} \langle 0|$$

$$H_{\text{out}} = |0\rangle_{\text{out}} \langle 0| \otimes |L\rangle_{\text{clock}} \langle L|$$

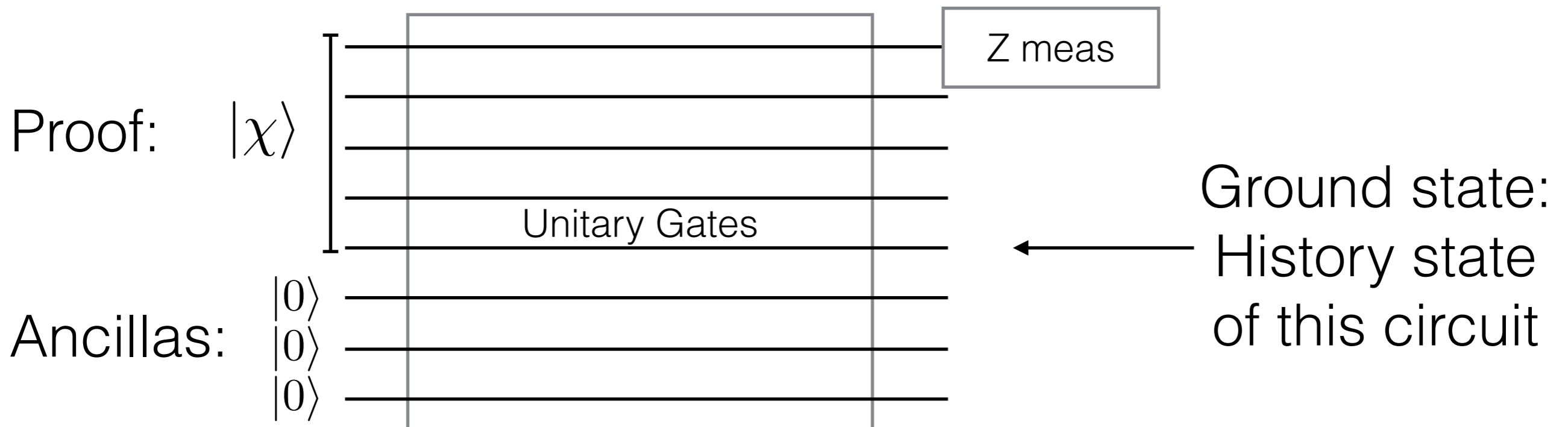
Encoding a computation in a ground state

Yes instance:

$$H = H_{\text{in}} + H_{\text{prop}} + H_{\text{out}}$$

No frustration:

Each term of Hamiltonian can reach its **ground space**:



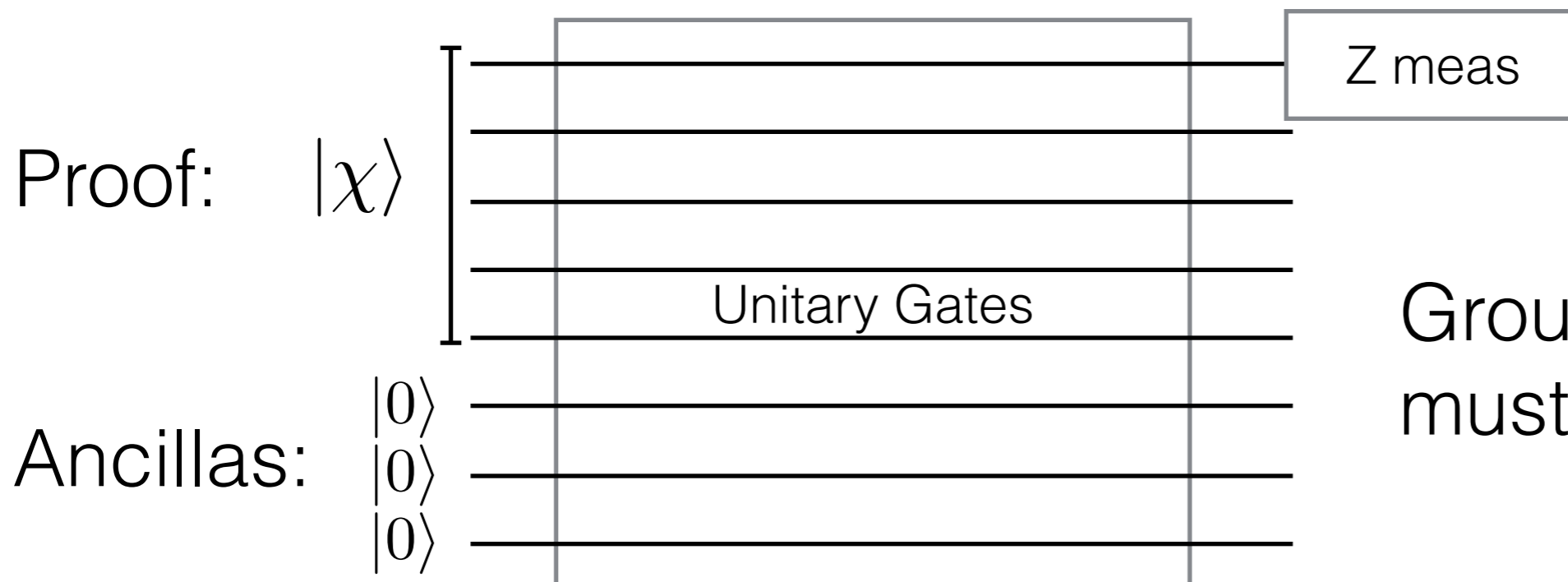
Encoding a computation in a ground state

No instance:

$$H = H_{\text{in}} + H_{\text{prop}} + H_{\text{out}}$$

Frustration:

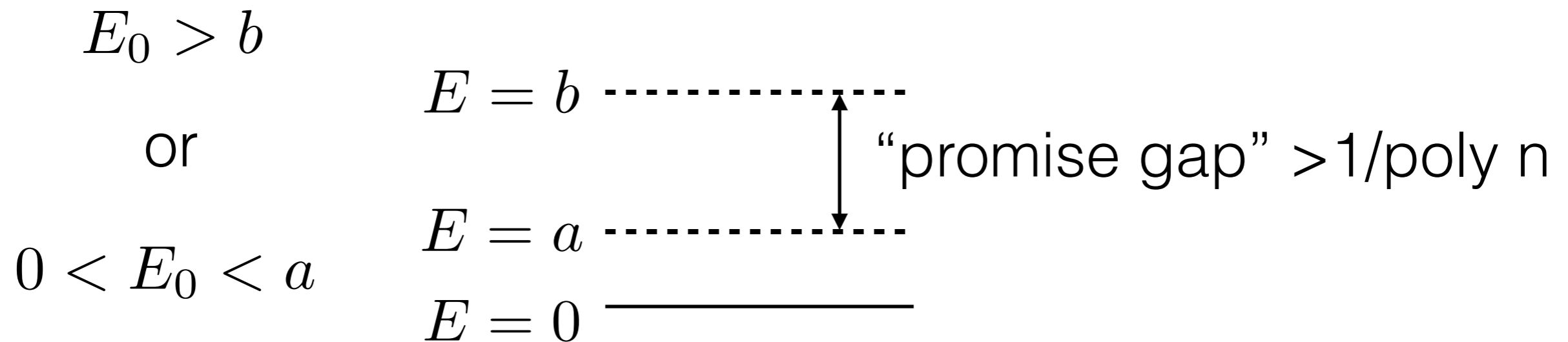
H_{out} wants output qubit to be in state 1, but no proof state exists to allow this.



Ground state energy must be **increased**.

Promise gap scaling

LH Problem: Determine whether $E_0 > b$ or $E_0 < a$?



We need the “promise gap” between **yes** and **no** to remain **inverse polynomial**.

Encoding a computation in a ground state

Kitaev:

Energy of **no** frustrated ground state instance scales with

$$E_0 \geq \lambda_1 2 \sin^2(\theta/2)$$

scales with $1/L$

2nd lowest eigenvalue of H_{prop}

Kitaev's H_{prop} :

Lowest eigenvalue:

$$1 - \cos(\pi/L + 1) \approx \frac{\pi}{2(L+1)^2}$$

LH is QMA-complete

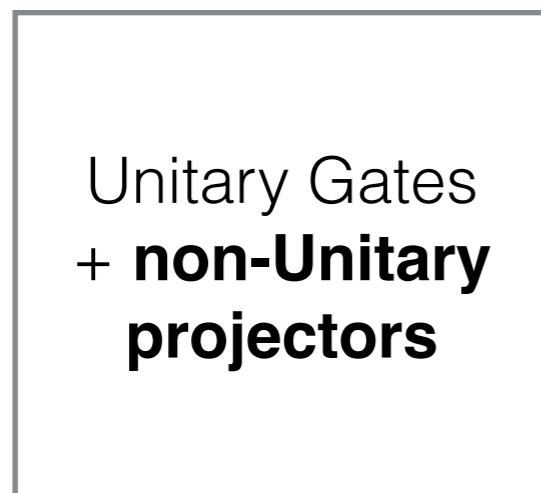
Non-unitary ground state computation?

Our work: (Usher/Browne, unfinished 2015)

Non-unitary ground state computation?

Our work: (Usher/Browne, unfinished 2015)

What if?



Ground State
Computation



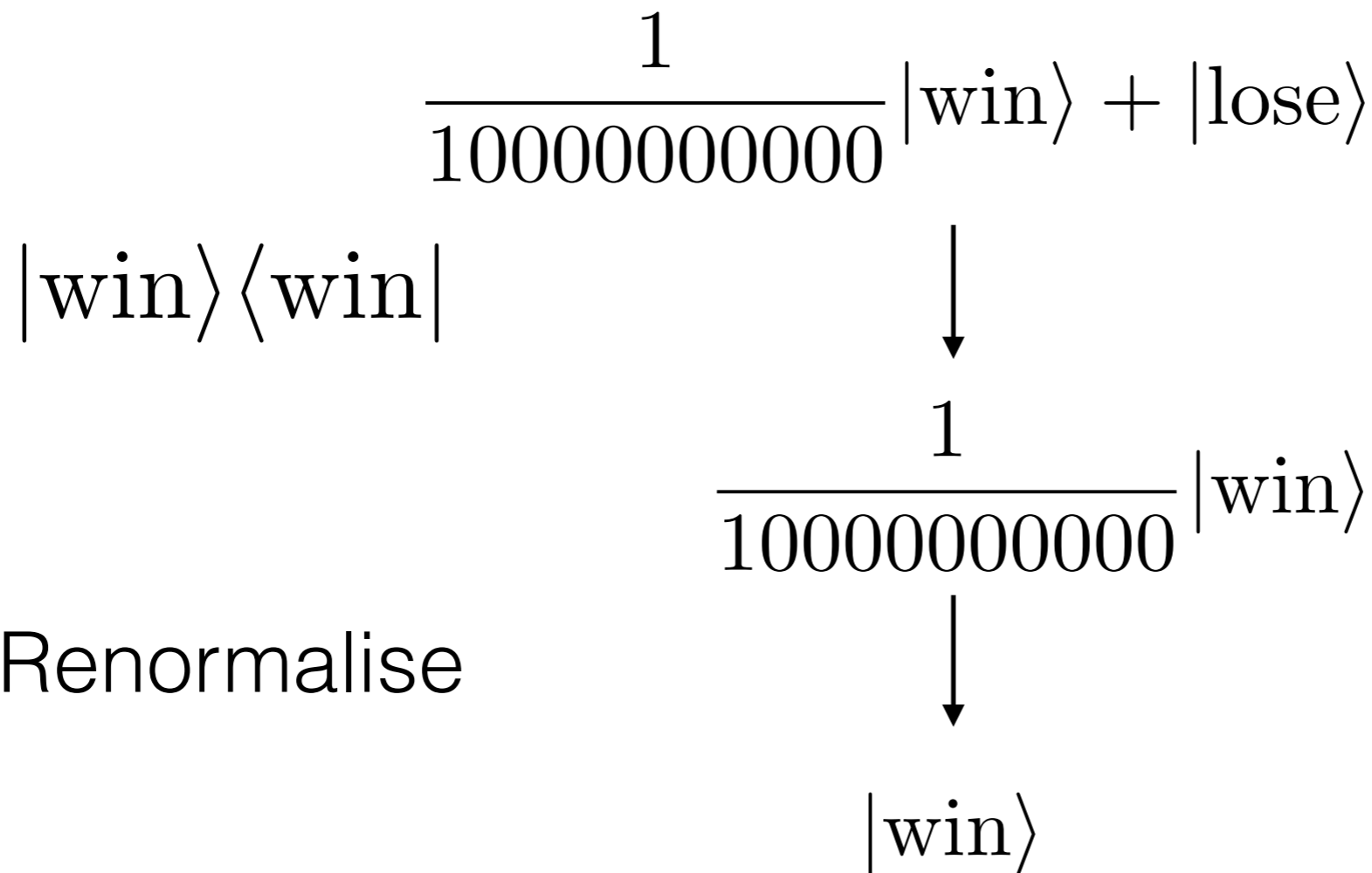
Add **projectors** to the
circuits

Verifier circuits

Non-unitary ground state computation?

Renormalised Projectors / Post-selection

E.g. quantum lottery ticket



Non-unitary ground state computation?

Why?

- Curiosity! - Feynman's construction pre-dates **all** quantum computing theory.
- Add **projectors** to unitary circuits and you get:

Fault-tolerant Quantum Computation

Measurement-based Quantum Computation

postIQP = postBQP = PP

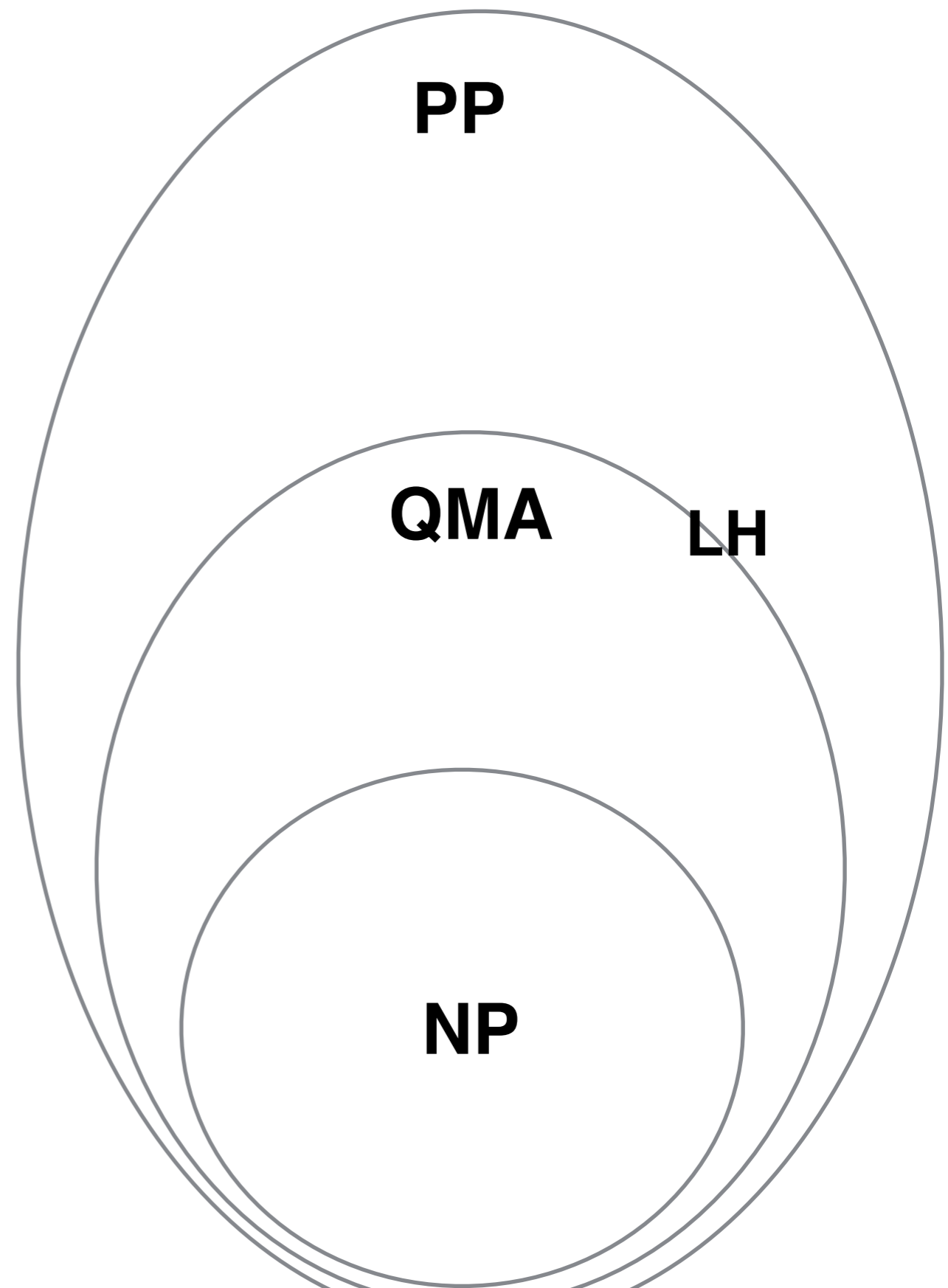
Non-unitary ground state computation?

Why it will never work!

$$\text{postIQP} = \mathbf{postBQP} = \mathbf{PP}$$

Aaronson:

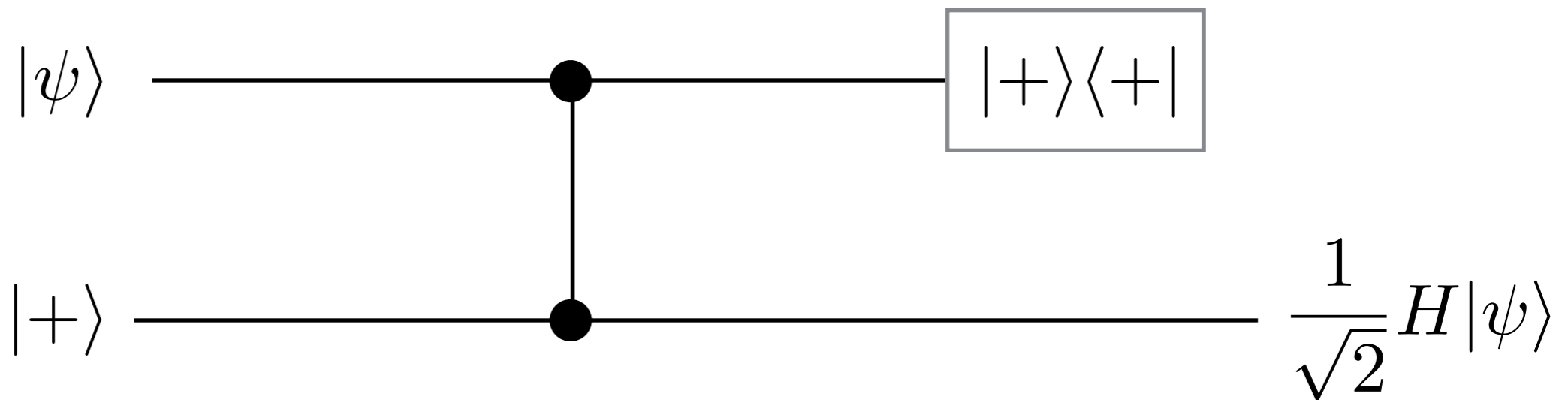
If deterministic projectors (post-selection) are added to unitary gates, we can efficiently solve **PP-hard** problems.



Non-unitary ground state computation?

Why it might just work.....

postlQP = postBQP = PP

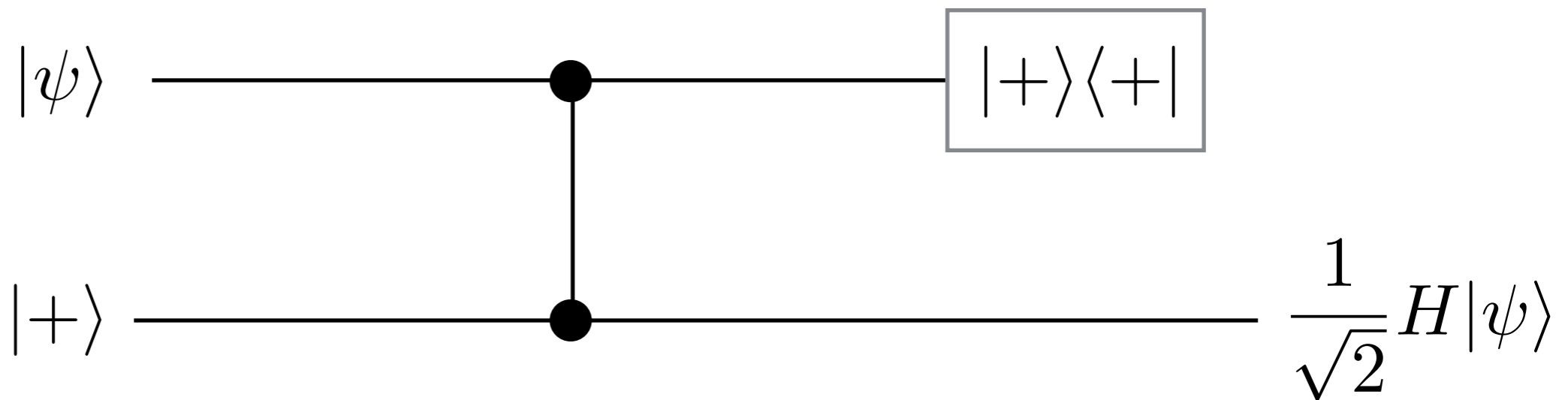


One-bit teleportation circuit (Zhou / Leung / Chuang 2000).

Non-unitary ground state computation?

Why it might just work.....

postlQP = postBQP = PP



One-bit Zero-bit teleportation circuit.

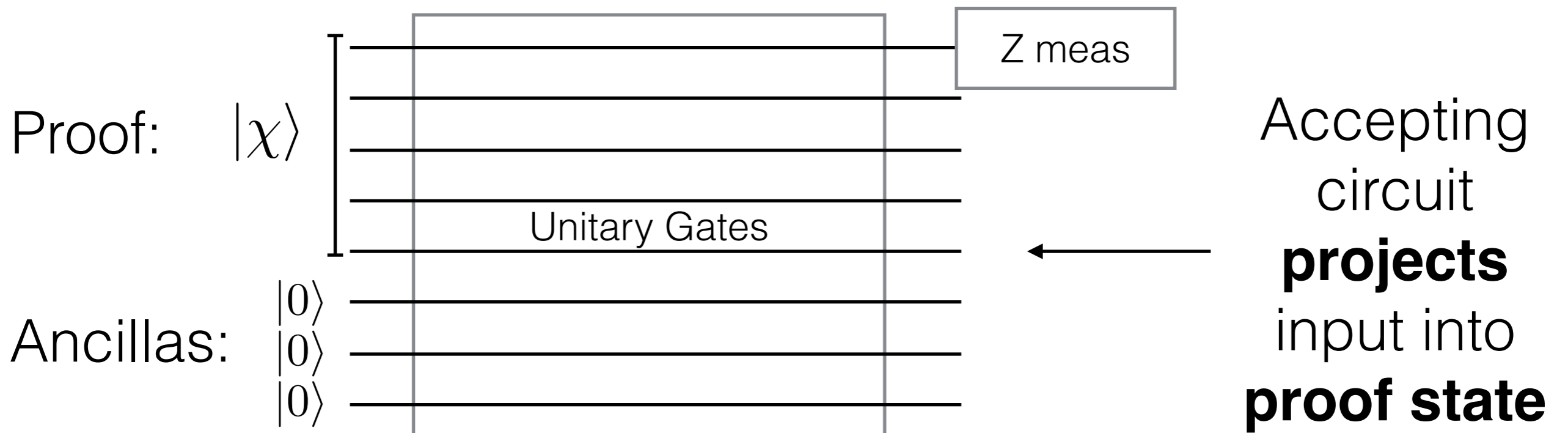
NB **Norm** of output is **independent** of input.

Encoding a computation in a ground state

Why it also might just work.....

In Kitaev construction, H_{out} **projects** the input to the correct proof state.

$$H = H_{\text{in}} + H_{\text{prop}} + H_{\text{out}}$$



The projector gadget

Recall Kitaev and Feynman's unitary gadget:

$$H_j = -\frac{1}{2} (U_j \otimes |j\rangle\langle j-1| + \text{h.c.}) + \frac{1}{2} \mathbb{1} \otimes (|j\rangle\langle j| + |j+1\rangle\langle j+1|)$$

that had ground space

$$|\psi\rangle|j-1\rangle + U_j|\psi\rangle|j\rangle$$

A projector gadget?

For projector P , want a gadget

$$H_p = ?$$

that has ground space

$$|\psi\rangle|j-1\rangle + \frac{P|\psi\rangle|j\rangle}{\beta}$$

where

$$\beta = \|P|\psi\rangle|j\rangle\|$$

To avoid PP-hardness, we **assume** β will be equal for **all** states $|\Psi\rangle$.

The projector gadget

$$H_j = \left(\frac{\beta^2}{1 + \beta^2} \right) P \otimes \left(-\frac{1}{\beta} |j-1\rangle\langle j| - \frac{1}{\beta} |j\rangle\langle j-1| + \frac{1}{\beta^2} |j-1\rangle\langle j-1| + |j\rangle\langle j| \right) + P^\perp \otimes |j\rangle\langle j|$$

has the ground space we want!

$$|\psi\rangle |j-1\rangle + \frac{P|\psi\rangle |j\rangle}{\beta}$$

NB In the limit $\beta \rightarrow 1$, $P \rightarrow \mathbb{1}$, we recover Kitaev/Feynman gadget.

The projector gadget

$$H_j = \left(\frac{\beta^2}{1 + \beta^2} \right) P \otimes \left(-\frac{1}{\beta} |j-1\rangle\langle j| - \frac{1}{\beta} |j\rangle\langle j-1| + \frac{1}{\beta^2} |j-1\rangle\langle j-1| + |j\rangle\langle j| \right) + P^\perp \otimes |j\rangle\langle j|$$

$$H_j \left(|\psi\rangle |j-1\rangle + \frac{P|\psi\rangle |j\rangle}{\beta} \right)$$

$$\propto -\frac{P}{\beta} |\psi\rangle |j\rangle + \frac{P}{\beta} |\psi\rangle |j\rangle = 0$$

The projector gadget

$$H_j = \left(\frac{\beta^2}{1 + \beta^2} \right) P \otimes \left(-\frac{1}{\beta} |j-1\rangle\langle j| - \frac{1}{\beta} |j\rangle\langle j-1| + \frac{1}{\beta^2} |j-1\rangle\langle j-1| + |j\rangle\langle j| \right) + P^\perp \otimes |j\rangle\langle j|$$

has the ground space we want!

$$|\psi\rangle |j-1\rangle + \frac{P|\psi\rangle |j\rangle}{\beta}$$

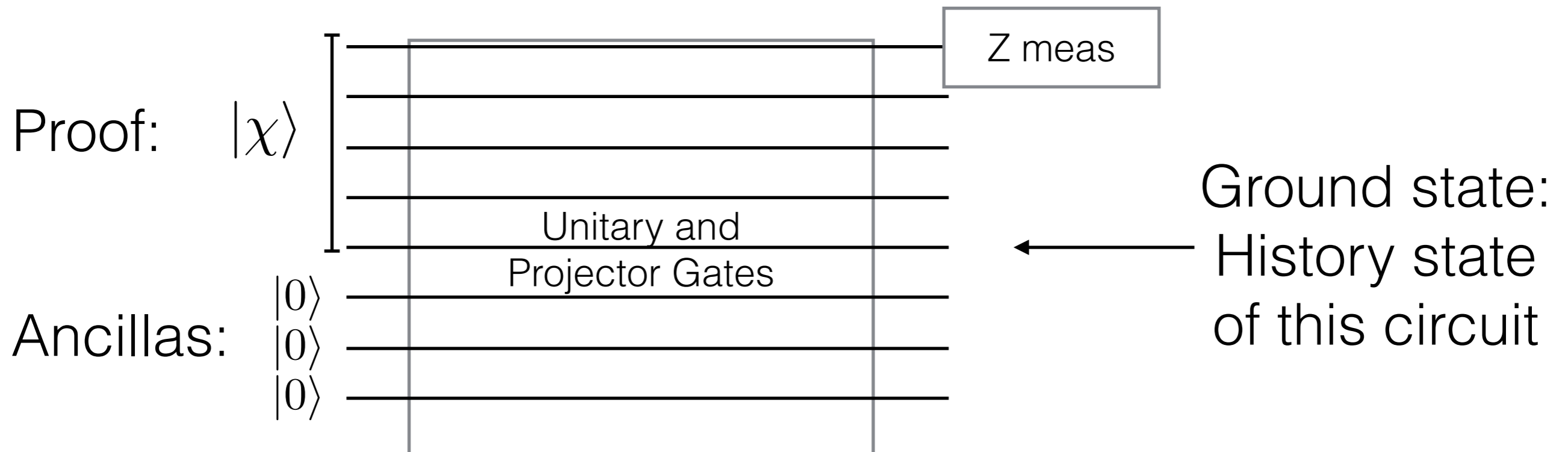
NB In the limit $\beta \rightarrow 1$, $P \rightarrow \mathbb{1}$, we recover Kitaev/Feynman gadget.

Encoding a computation in a ground state

With this gadget, we can construct Hamiltonians

$$H = H_{\text{in}} + H_{\text{prop}} + H_{\text{out}}$$

such that “**Yes**” instances have low energy...



Encoding a computation in a ground state

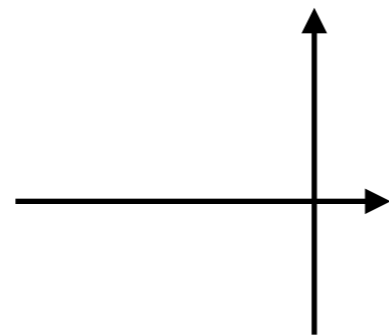
And **No** instances have frustration

$$H = H_{\text{in}} + H_{\text{prop}} + H_{\text{out}}$$

The “promise gap” satisfies Kitaev’s formula

$$E_0 \geq \lambda_1 2 \sin^2(\theta/2)$$

Need to check this!



should **scale** as **1 / poly(n)**

2nd lowest eigenvalue of H_{prop}

Characterising the promise gap

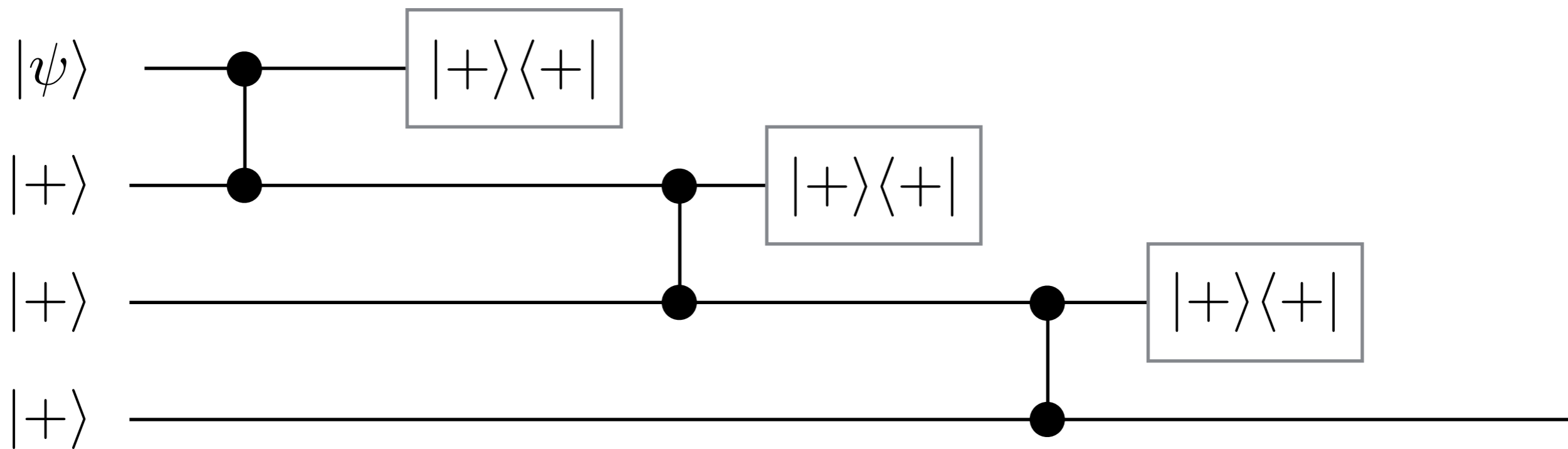
December 2014: We had a beautiful and elegant **analytic bound** on λ_1 for many circuits.

January 2015: We found one of **these** in the proof...

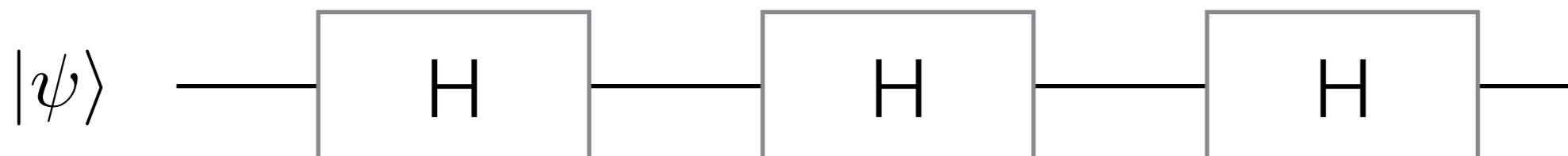


Characterising the promise gap

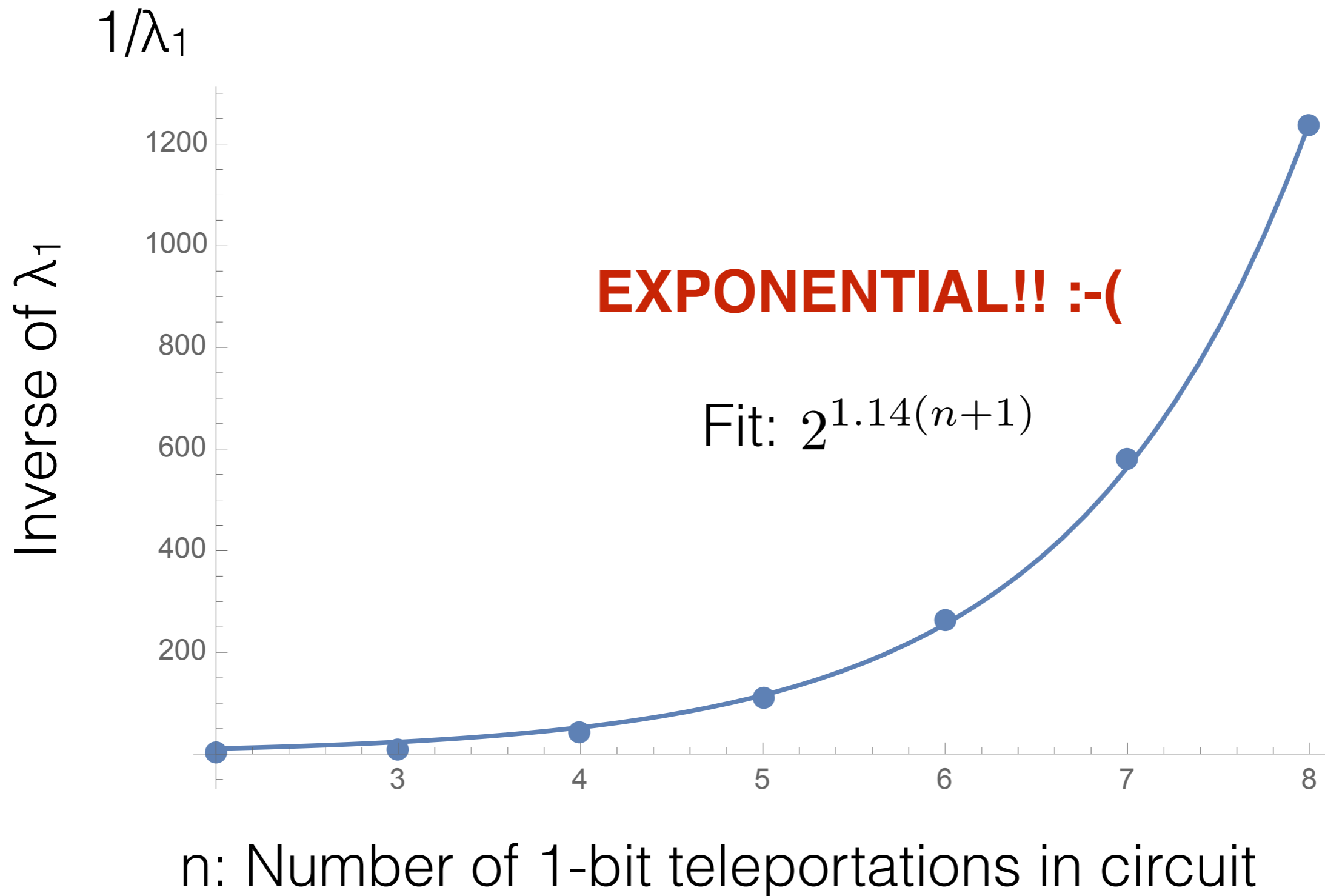
Instead - some quick and simple proof-of-principle numerics.



Equivalent to

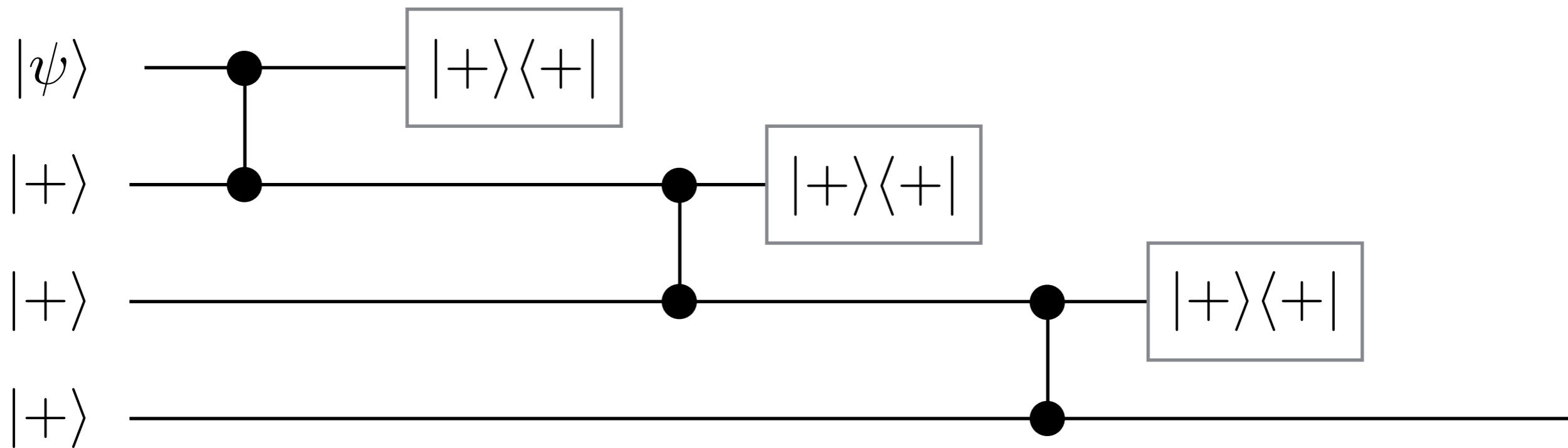


Characterising the promise gap



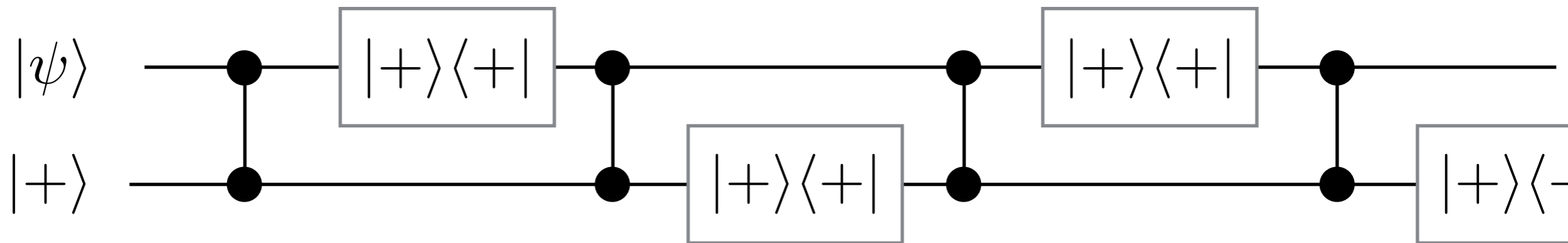
Characterising the promise gap

Extra ancillas giving **low energy** excited states. :-)

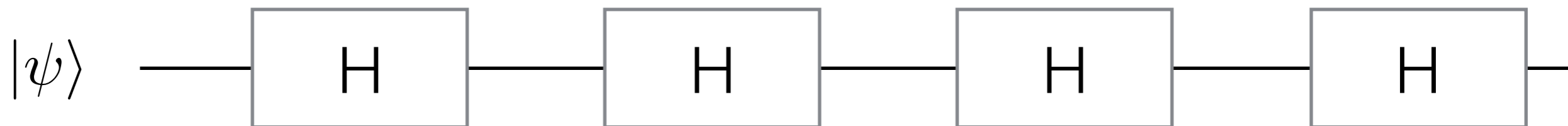


Characterising the promise gap

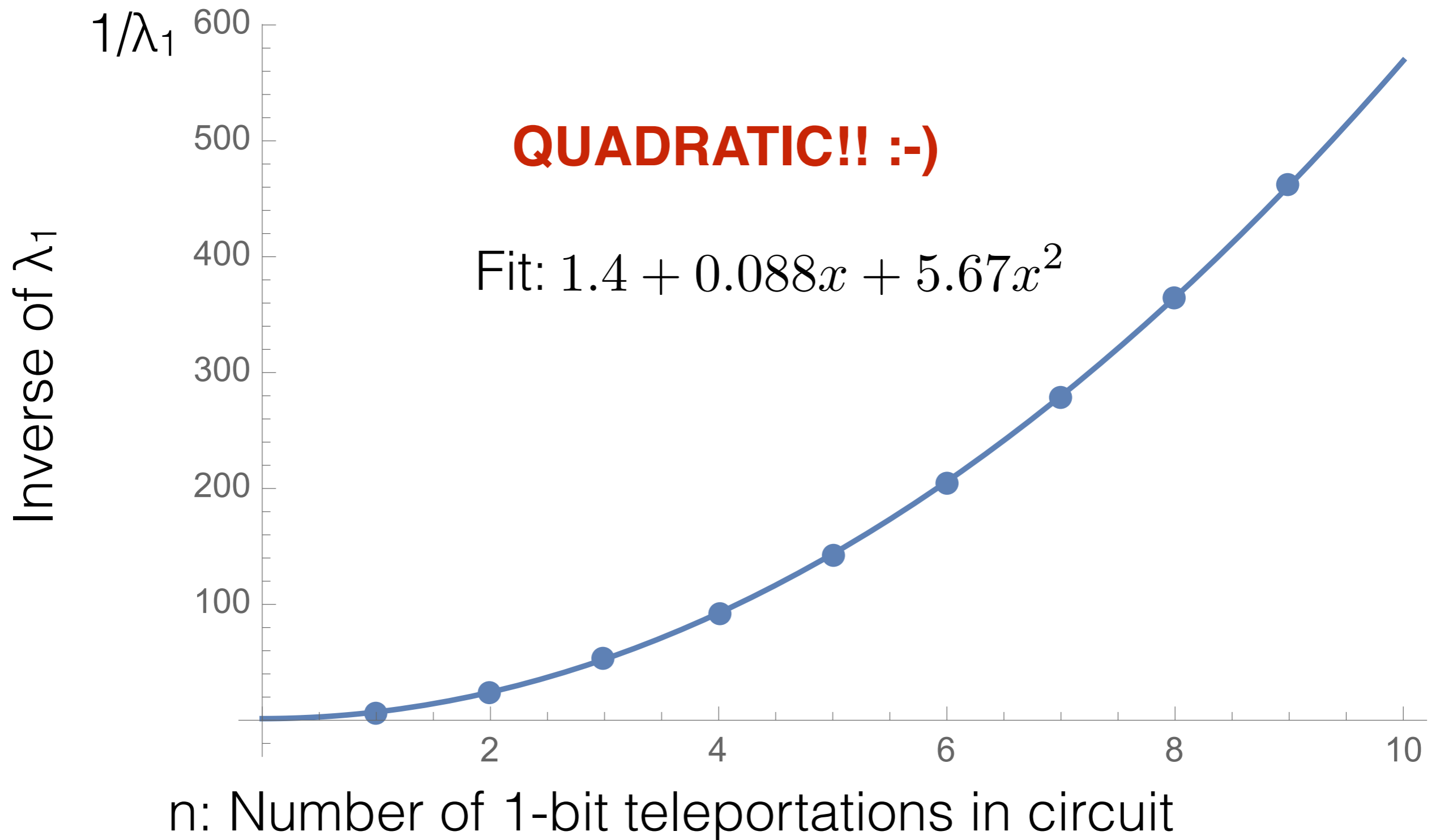
A second try



Equivalent to



Characterising the promise gap



Conclusion

It **works**! Sometimes.....

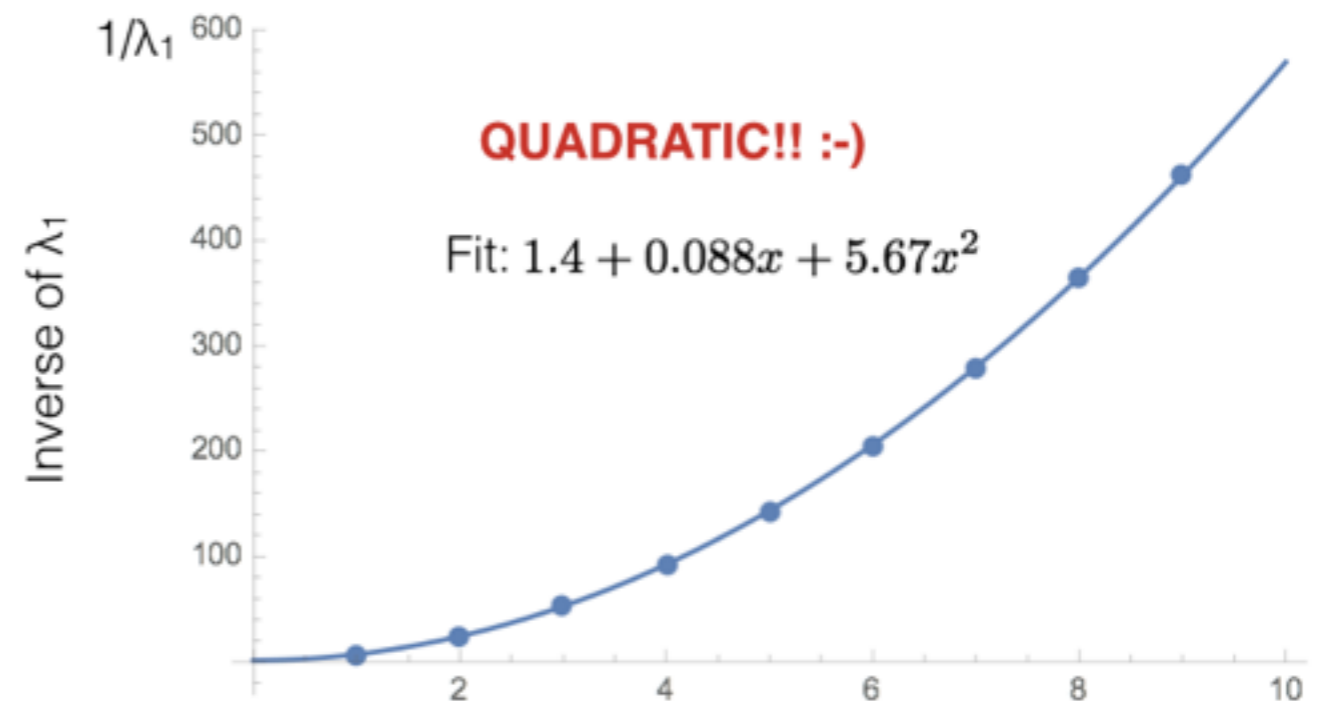
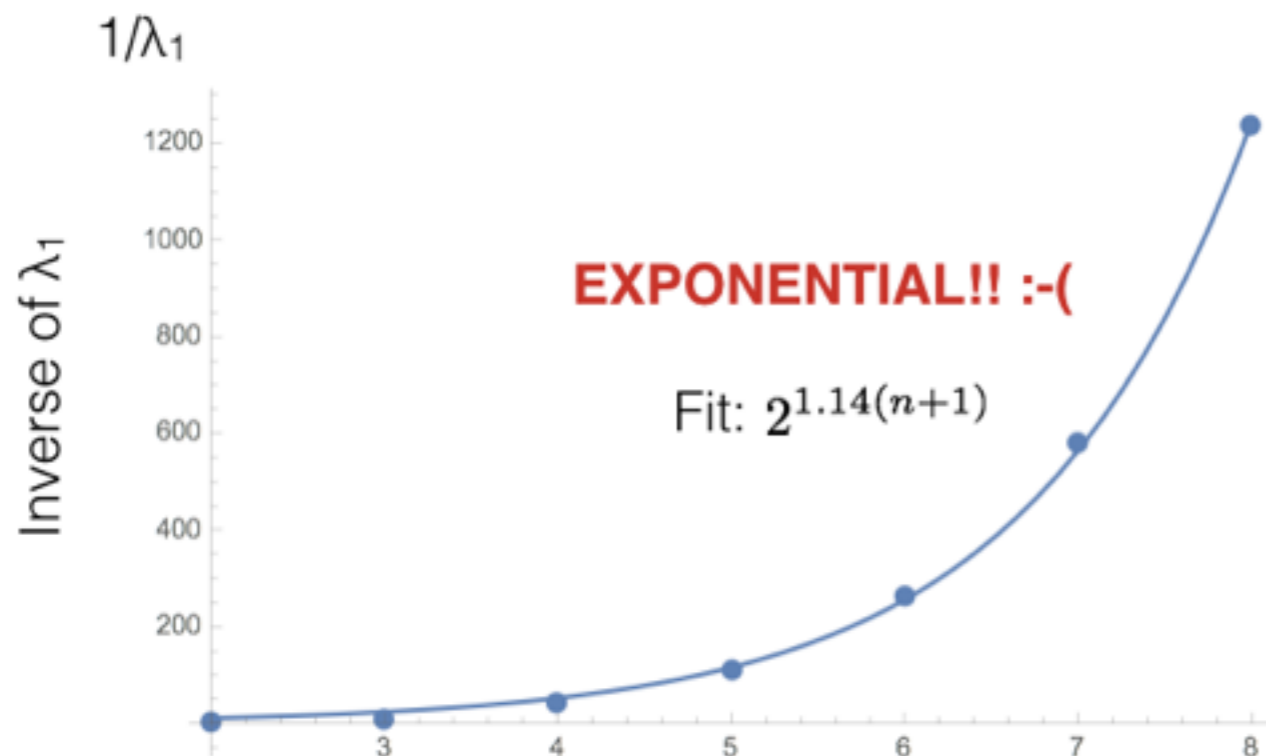
(probably)

Where do we go next?



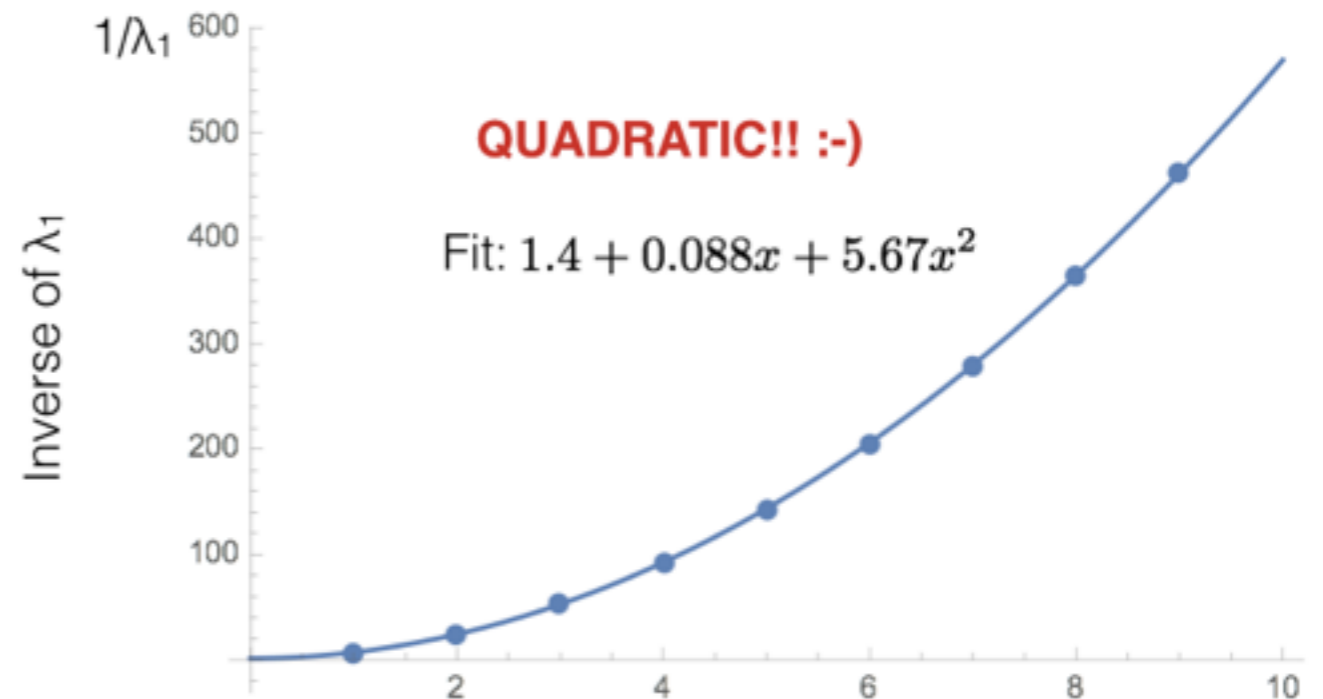
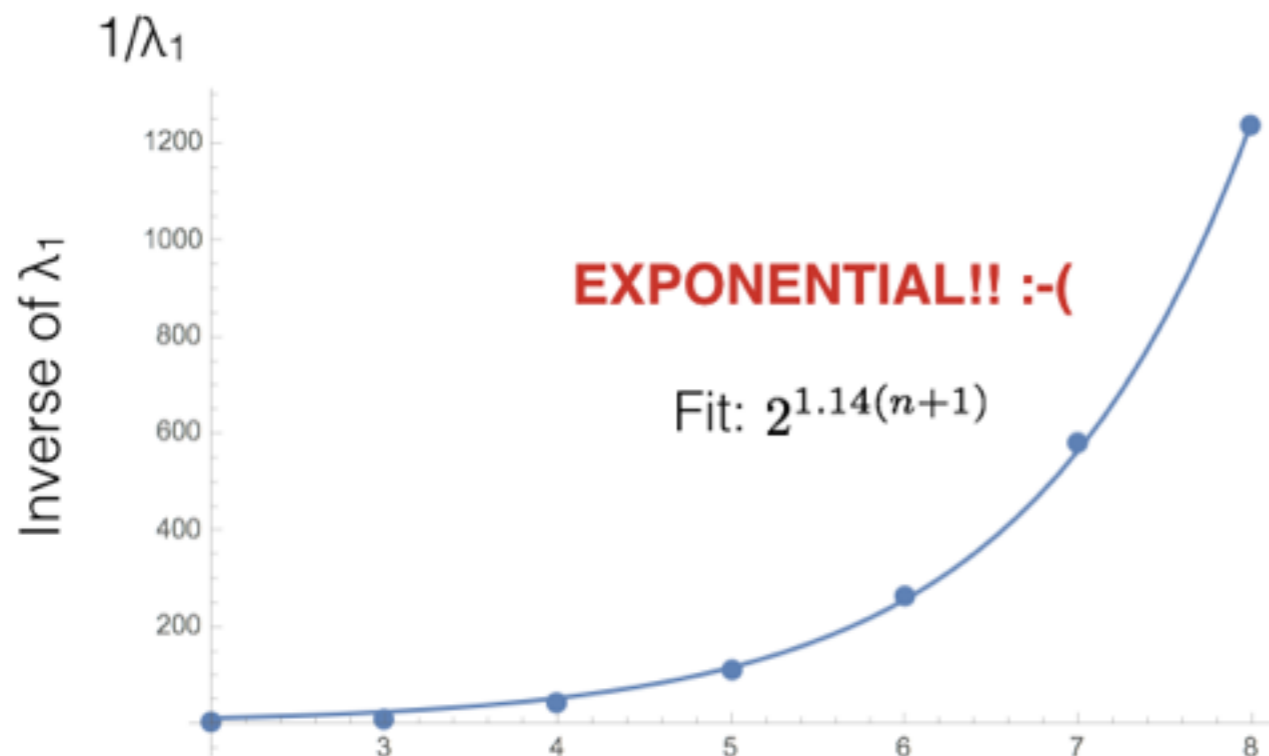
Where do we go next?

- ❖ Numerical evidence for 1/poly promise gap scaling:
 - ❖ But when does **exponential scaling** occur?
 - ❖ Analytic bounds?
- ❖ Aim: A QMA construction from a **non-universal** gate set with projectors (e.g. **IQP** circuit).



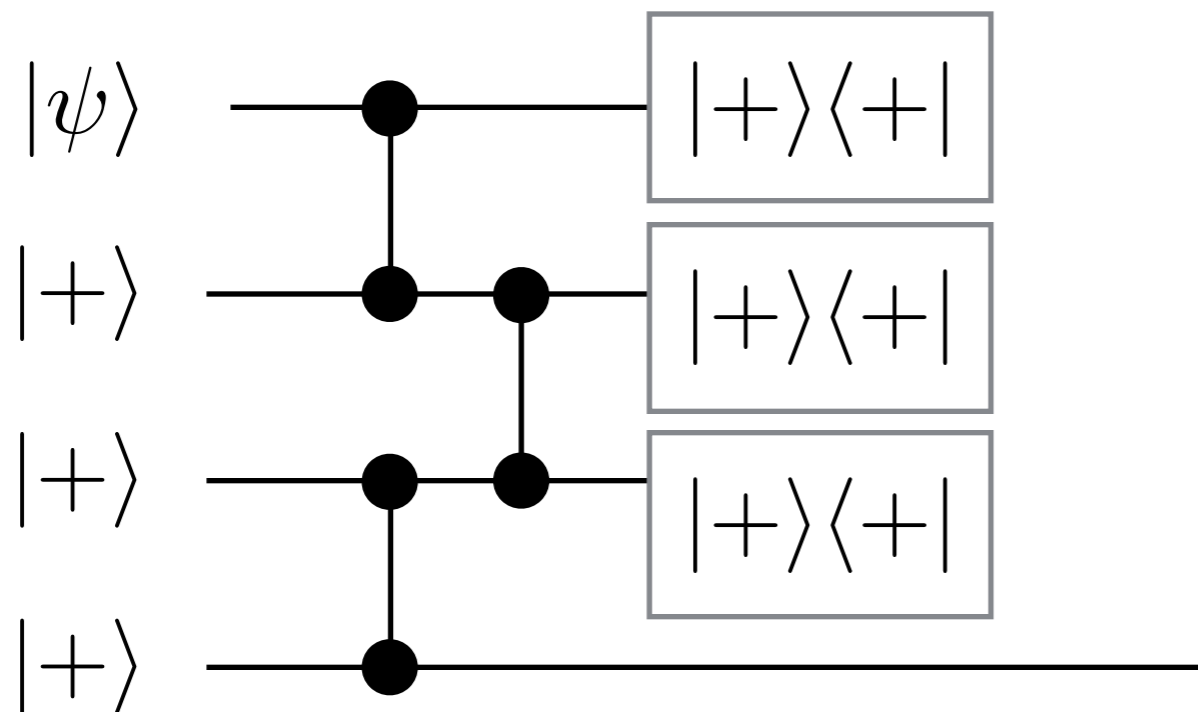
Where do we go next?

- ❖ Complexity
 - ❖ When does **BQP + projectors = BQP**
 - ❖ Is **constant probability** on **input states** sufficient?
 - ❖ Can we characterise **exponentially closing gap** with other complexity classes?



Where do we go next?

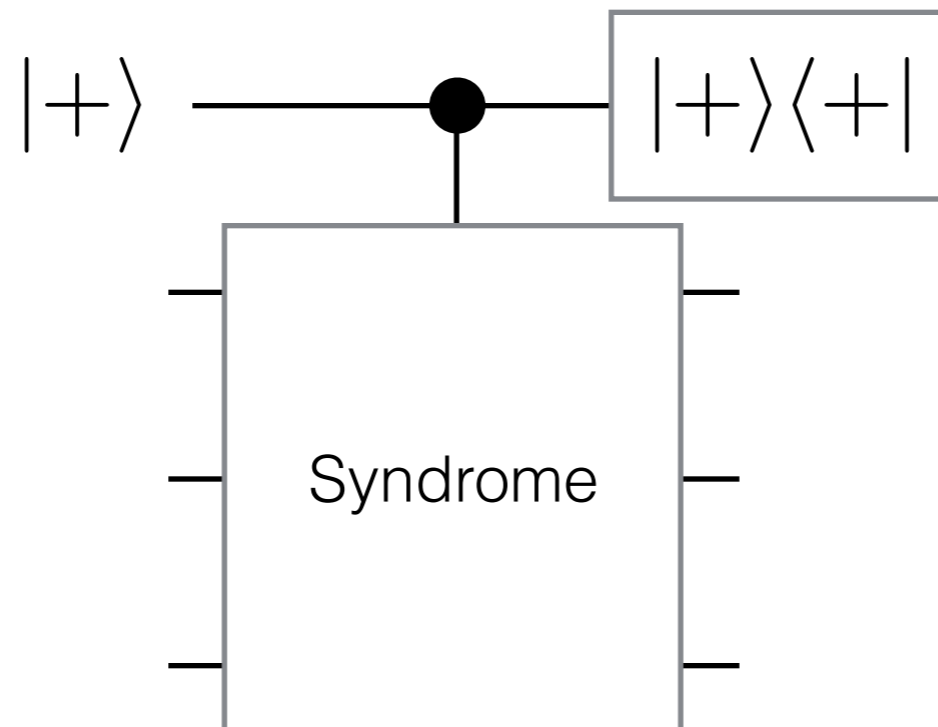
- ❖ Post-selection gives **IQP** / **MBQC** circuits trivial time complexity. QMA with **constant clock-steps**?



Need **many** clocks!

Where do we go next?

- ❖ Incorporate **fault tolerance**?
- ❖ **Error detection** gadgets to project onto **error free** states?
- ❖ Robustness of Hamiltonian to **perturbations**?
- ❖ Norm of history state vectors problematic?



Where do we go next?

- ❖ **Other applications** for the projector gadget?
- ❖ **Adiabatic** Quantum Computation?
- ❖ Relationship with Bacon and Flammia's **adiabatic cluster state** model?

Adiabatic One-way QC



$$H = -Z_{n-1}X_n - \sum_{j=2}^{n-1} S_j \quad S_j = Z_{j-1}X_jZ_{j+1}$$

Thank you!

