

# L13 Ferromagnetism

## Lecture outline:

- The mechanism of ferromagnetism
- Domains.
- Antiferromagnetism.
- Ampere s law in a magnetic material.
- Electromagnets
- Magnetic recording

# L13.1 Ferromagnetism

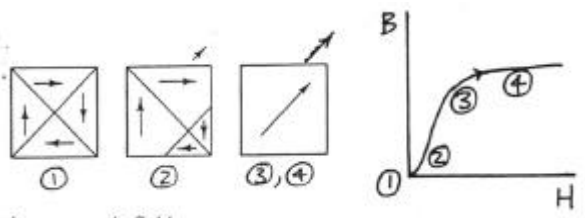
## The Mechanism of Ferromagnetism

Iron crystals are made of domains, in which the magnetic dipoles are aligned.

4 domains:

(1) Dipoles cancel,  
not magnetized

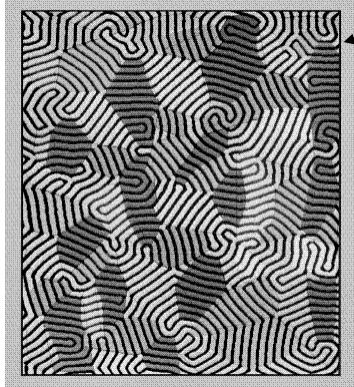
(2) Domains grow



(3) Some domains  
disappear,

(4) All dipoles  
aligned, saturation

## L13.2 Ferromagnetism



Magnetic domains as seen in a polarizing microscope

There is a critical temperature, the Curie temperature, above which the thermal motion randomizes the domains and destroys the magnetization.

$T_c$

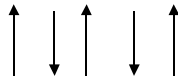
Fe 1043 K

Co 1388 K

Ni 627 K

## L13.3 Ferromagnetism

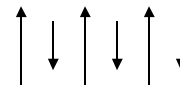
Anti-ferromagnetism: interaction between neighbouring dipoles has opposite sign: adjacent moments tend to align antiparallel ( paramagnetic ), eg Cr, Mn.



Ferrimagnetism:

If there are 2 unequal kinds of moments, there is a net magnetization even under complete antiparallel ordering:

Such materials are ferrites, eg lodestone,  $\text{FeOFe}_2\text{O}_3$ .



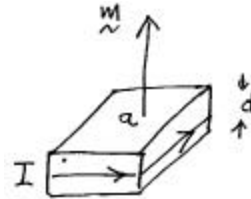
## L13.4 Ferromagnetism

Consider a little loop: dipole moment  $\mathbf{m}$

Magnetization  $\mathbf{M} = \mathbf{m}/\text{volume} = \frac{\mathbf{m}}{ad}$

$m = Mad = Ia$  so  $I = Md$

and  $\frac{I}{d} =$  "surface current density" ( $\text{Am}^{-1}$ )  $= M = J_{\text{surface}}$



(compare result relating polarization to surface charge density)

## L13.5 Ferromagnetism

Ampere's law:

$$\oint \mathbf{B} \cdot d\mathbf{s} = \mu_0 \oint (\mathbf{H} + \mathbf{M}) \cdot d\mathbf{s} = \mu_0 i = \mu_0(i_{\text{free}} + i_{\text{surface}})$$

Here  $i_{\text{surface}} = i_{\text{bound}}$  (due to internal dipoles)

$i_{\text{free}}$  = current applied from outside



$$\uparrow \mathbf{M} = 0$$

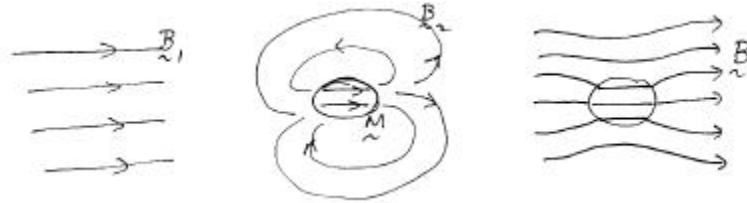
But  $\mu_0 \oint \mathbf{M} \cdot d\mathbf{s} = \mu_0 I = \mu_0 i_{\text{surface}}$

so  $\oint \mathbf{H} \cdot d\mathbf{s} = i_{\text{free}}$

This is Ampere's law in a magnetic material.

## L13.6 Ferromagnetism

Put a magnetized sphere in a magnetic field:

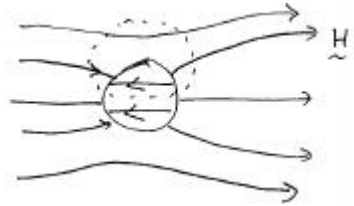


**M** in sphere is uniform. **B** is continuous (can't start or stop).

But **H** may not be continuous:

On the path,  $\oint \mathbf{H} \cdot d\mathbf{s} = 0$

so **H** may reverse in the material.

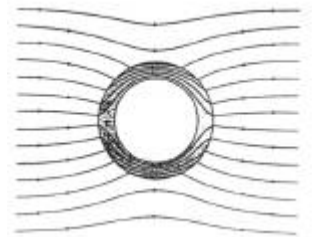


School of Physics- N. Cramer and R. McPhedran 2001

7

## L13.7 Ferromagnetism

Magnetic shielding: put a hole in a ferromagnet. The field gets trapped in the magnet, with zero field inside hole. (compare electrical shielding).



School of Physics- N. Cramer and R. McPhedran 2001

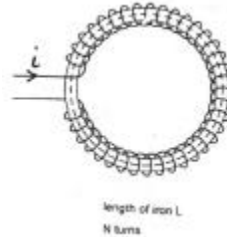
8

## L13.8 Ferromagnetism

Use Ampere's law in an iron toroid:

$$\oint \mathbf{H} \cdot d\mathbf{s} = i_{free} = Ni = HL$$

so  $H = \frac{Ni}{L}$  and  $B = \mu \frac{Ni}{L}$  inside.



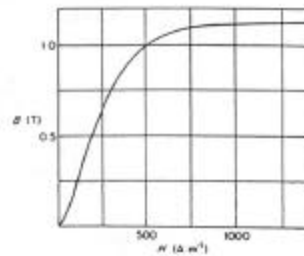
Example:  $Ni = 1000$ ,  $L = 1$  m,

so  $H = 1000 \text{ Am}^{-1}$ .

From the chart,  $B = 1.15 \text{ T}$ ,

so  $\mu = B/H = 1.15 \times 10^{-3}$ ,

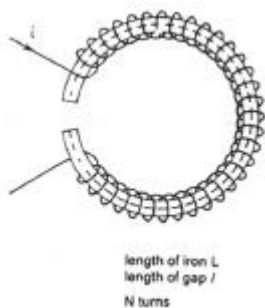
and  $\mu_r = \mu/\mu_0 = 1.15 \times 10^{-3}/(4\pi \times 10^{-7})$   
 $= 900$ .



School of Physics- N. Cramer and R. McPhedran 2001

## L13 Ferromagnetism

### Electromagnet



with a gap. For a small gap,  $\mathbf{B}$  is continuous, so  $\mathbf{B}$  is nearly the same in the iron and the gap.

Now  $H$  in the iron is  $H_1 = \frac{B}{\mu} = \frac{B}{\mu_r \mu_0}$

and  $H$  in the gap is  $H_2 = \frac{B}{\mu_0}$

so  $\oint \mathbf{H} \cdot d\mathbf{s} = H_1 L + H_2 l = \frac{BL}{\mu_r \mu_0} + \frac{Bl}{\mu_0} = Ni$

$\therefore B = \frac{\mu_0 Ni}{L/\mu_r + l} \approx \frac{\mu_0 Ni}{l}$  because  $\mu_r$  is very big.

Example:  $L = 1$  m,  $l = 1$  cm,  $Ni = 1000$ , then  $B = 4\pi \times 10^{-7} \times 1000/0.01$   
 $= 0.13 \text{ T}$

School of Physics- N. Cramer and R. McPhedran 2001

## L13.10 Ferromagnetism

For a permanent magnet,  $i = 0$ , so

$$H_1 L + H_2 l = 0 \quad \text{so} \quad H_{iron} = -\frac{H_{gap} l}{L} = -\frac{Bl}{\mu_0 L}$$

(note the minus sign)



Example:  $L = 1\text{m}$ ,  $l = 1\text{cm}$ ,  $B = 0.1\text{T}$ ,

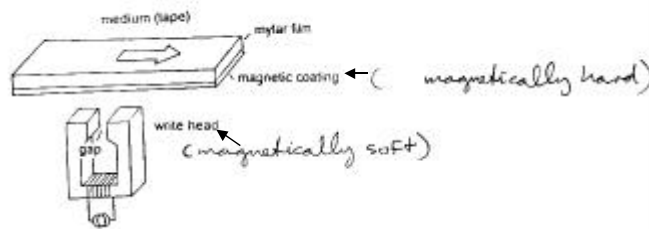
$$H_{gap} = \frac{B}{\mu_0} = \frac{0.1}{4\pi \times 10^{-7}} \simeq 80000\text{Am}^{-1} \quad \text{and} \quad H_{iron} = -800\text{Am}^{-1}$$

School of Physics - N. Cramer and R. McPhedran 2001

11

## L13.11 Ferromagnetism

Magnetic recording:



There are 3 heads: erase, write and read. The tape with magnetic coating moves over the heads. At the gap in the electromagnet, there is a fringing magnetic field that extends over the tape.



In writing, the time variation in the field is translated to a space varying magnetization in the moving tape.

edran 2001

12

## L13.12 Ferromagnetism

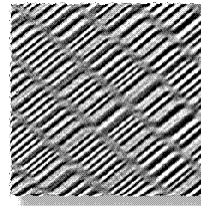
Audio tape uses needle-like particles of  $\text{Fe}_2\text{O}_3$  or  $\text{CrO}_2$  ( $\sim 1\mu\text{m}$ ), on a substrate of mylar (a polyester).



The spatial frequency translates to an audio frequency (analogue).  
Video tape is similar.

In reading, the changing field in the tape induces a signal in the coil.

Computer disks store information  
digitally, as bits. Domains are magnetized in up or down direction.



Field of view  
 $30\mu\text{m}$ .