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Quantum Mechanics

Lecture 13

Identical particles;
Spin-statistics theorem;
Helium atom;
Exchange interaction.



A quick recap

Perturbation theory can be used to estimate energies and eigenstates when the complete Hamiltonian is too complicated to solve explicitly.

Examples:

$$H_{SO} \propto \mathbf{S} \cdot \mathbf{L}$$

Spin-orbit coupling

$$H_{HF} \propto \mathbf{S} \cdot \mathbf{I}$$

Hyperfine structure

$$H_{\text{Stark}} \propto e |\mathbf{E}|$$

Stark shift

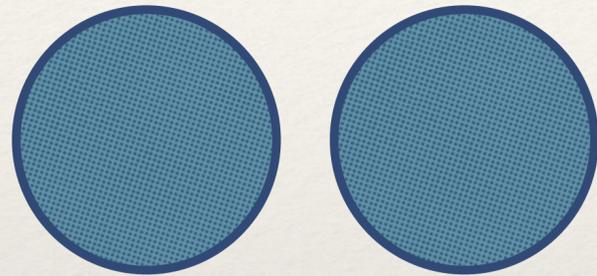
In some cases, the “old” quantum numbers for AM become “bad” and must be replaced by new quantum numbers for total AM. The Clebsch-Gordan coefficients tell us how to express the new eigenstates in terms of the old.

$$|JM\rangle = \sum_{m_1=-j_1}^{j_1} \sum_{m_2=-j_2}^{j_2} C_{m_1 m_2 M}^{j_1 j_2 J} |j_1 m_1 j_2 m_2\rangle$$

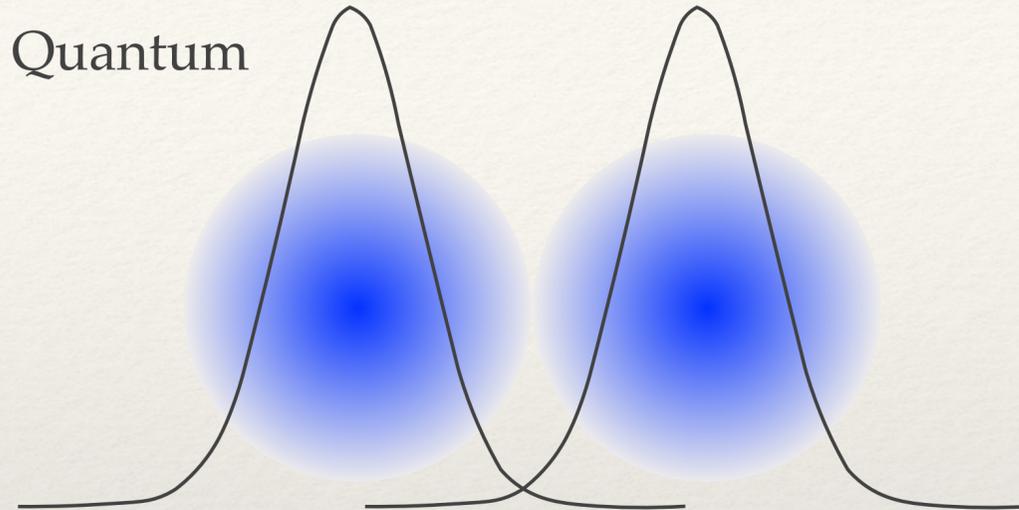
Identical particles

Elementary particles of the same type are all **exactly identical**.

Classical



Quantum



When the wave functions of identical particles overlap, they become **indistinguishable**. Contrast this with classical mechanics, where particles always occupy distinct space and can always (in principle) be distinguished.

This indistinguishability has important physical consequences.

Identical particles

Consider an experiment with two **indistinguishable** particles where we measure position in small regions dr . We must have:

Introduce the **exchange operator** (also known as **swap**):

Thus, indistinguishable particles (invariant under exchange) come in two types:

Bosons

Fermions

Bosons, Fermions, and the spin-statistics theorem

Bosons have integer spin, and fermions have half-integer spin.

In relativistic quantum mechanics, this is a provable statement known as the **spin-statistics theorem**.

The exchange phases are internally consistent under composition of particles:

Standard Model of Elementary Particles

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
	u up	c charm	t top	g gluon	H higgs
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	

QUARKS (purple text on the left side of the quark section)

LEPTONS (green text on the left side of the lepton section)

GAUGE BOSONS VECTOR BOSONS (red text on the right side of the gauge boson section)

SCALAR BOSONS (yellow text on the right side of the scalar boson section)

Symmetrizing wave functions

When quantum statistics are important, we need to explicitly enforce the exchange symmetry in our wave functions.

Consider non-interacting fermions in a spin-independent potential:

We can make this anti-symmetric with the use of the **Slater determinant**:

Ex: $N = 3$

The Pauli exclusion principle: this vanishes whenever $\psi_i = \psi_j$.

Symmetrizing wave functions

If the Hamiltonian is spin-independent and non-interacting, the wave function is a product of spatial and spin degrees of freedom.

Example: 2 particles

Symmetrization of spin has different consequences for bosons and fermions.

Example: 2 particles

spin-0 spin states:

spin-1/2 spin states:

Symmetrizing wave functions

The **total** wave function must have appropriate symmetry.

Example: 2 particles

In the non-interacting case, we can symmetrize space and spin separately.

spin-0 Bosons:

spin-1 / 2 Fermions:

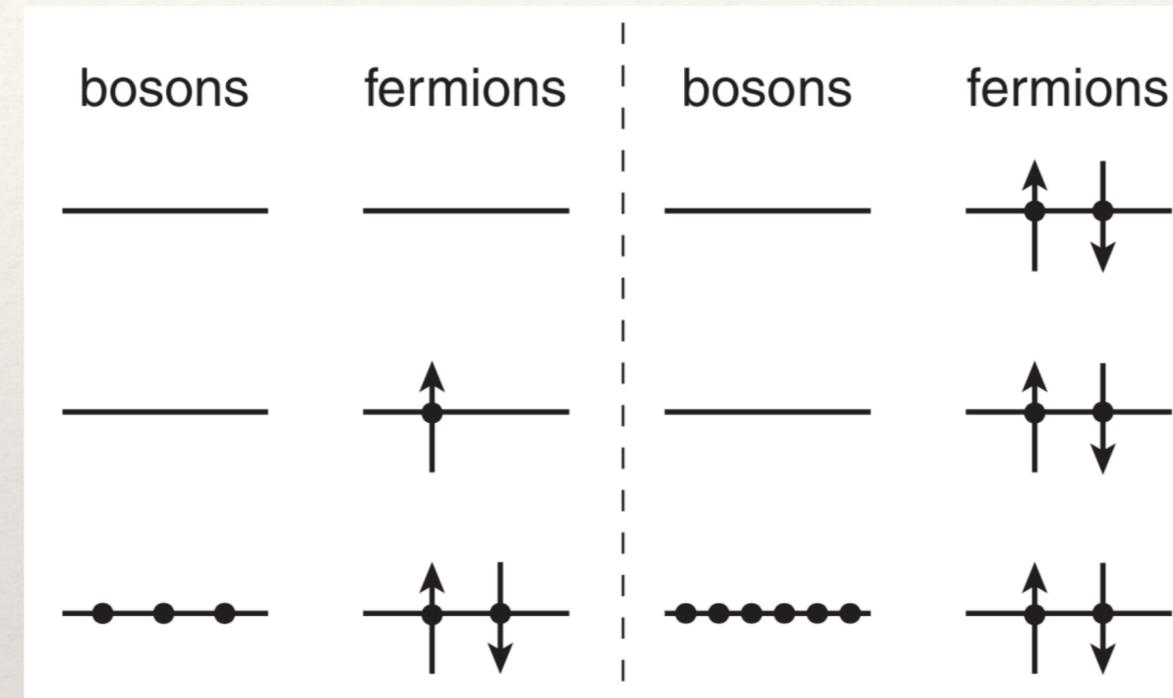
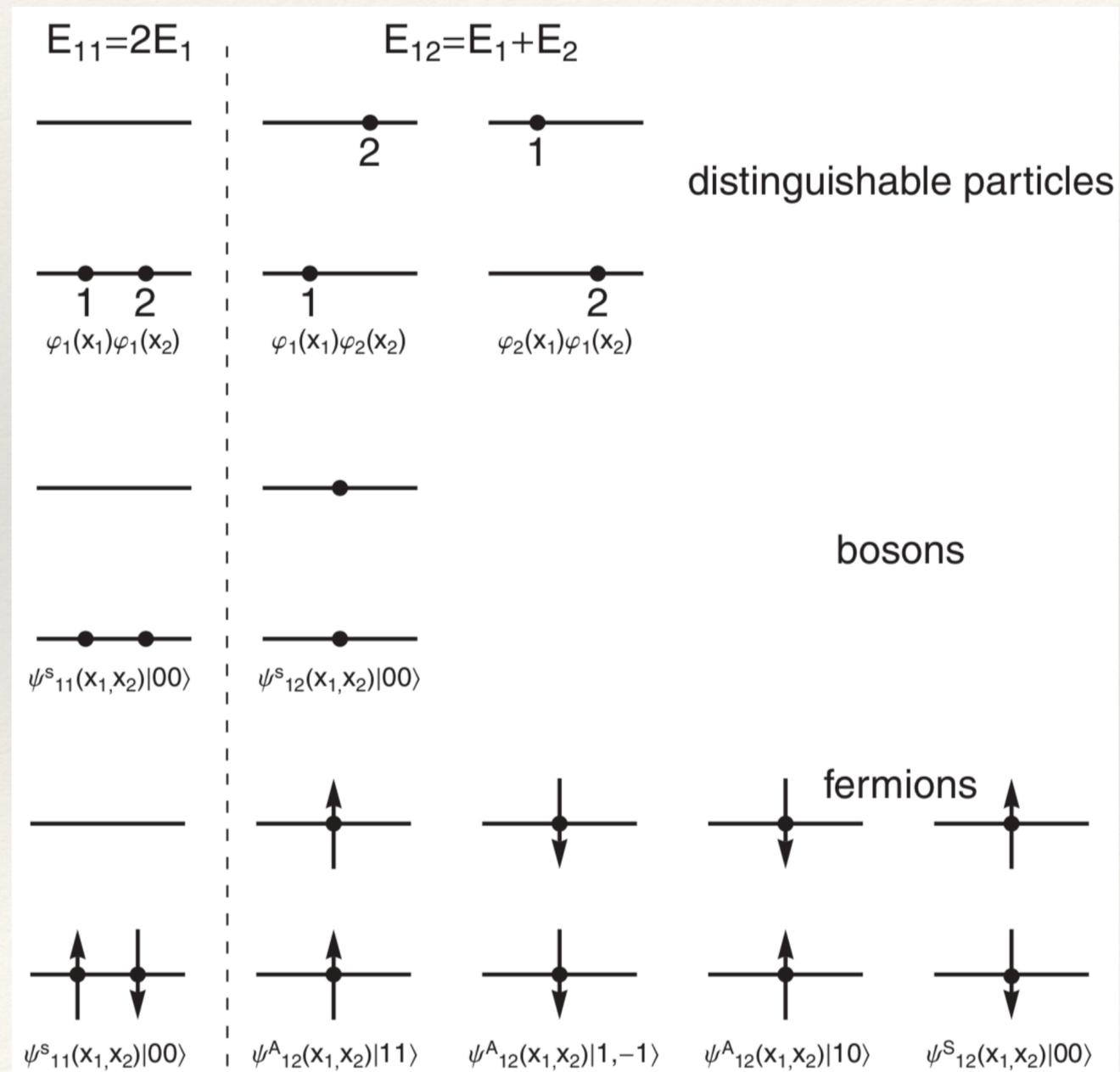
Ground
state:

Excited
states:

Fermionic excited states are four-fold degenerate, other states are unique.

Particle statistics

Particle symmetry leads to different degeneracies, and eventually to different notions of particle statistics.



These differences eventually lead to **Bose-Einstein statistics** for bosons and to **Fermi-Dirac statistics** for fermions.

Helium atom

The Hamiltonian for Helium, ignoring spin-orbit coupling and other spin-dependent effects is:

It is natural to try perturbation theory starting with $H_1 + H_2$ eigenstates.

Ground state:

Helium atom

There are several excited states with the requisite symmetry:

The explicit energy integrals in position space only depend on the spatial part:

Energy shift in He

The energy shift depends on l and the symmetry of the spatial wave function

$$E^{(1)}(l = 0) = 11.4\text{eV} \pm 1.2\text{eV} \quad E^{(1)}(l = 1) = 13.2\text{eV} \pm 0.9\text{eV}$$

Spin-dependent energy shifts have appeared from spin-free Hamiltonian interactions!

This effect is called the **exchange interaction**.

Antisymmetric wave functions tend to avoid overlapping in space, so they are farther apart and have less Coulomb repulsion.

Symmetric wave functions clump together and therefore have larger Coulomb repulsion.

