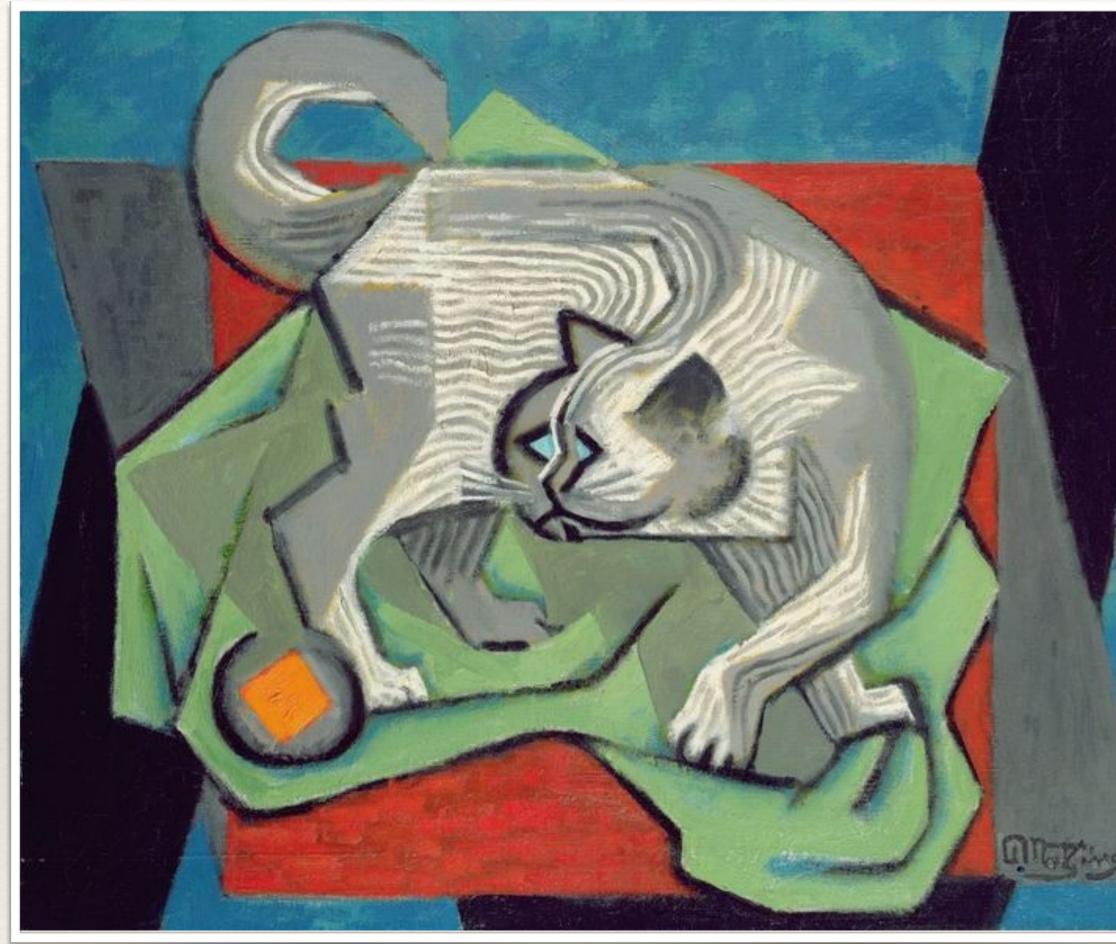
Prof. Steven Flammia

Quantum Mechanics

Lecture 5

Quiz 1; Eigenvalue spectrum; Angular momentum quantum numbers; Matrix elements.





A quick recap

Rotation operators are generated by angular momentum (AM): $R(\theta \mathbf{n}) = e^{i\theta \mathbf{n} \cdot \mathbf{J}/\hbar}$

AM operators obey the relations:

$$[J_i, J_j] = i\hbar \sum_k \epsilon_{ijk} J_k \qquad J^2 = J_x^2 + J_y^2 + J_z^2$$

Define raising and lowering operators that obey:

$$J_{\pm} = J_x \pm i J_y \qquad \qquad J_{\pm}^{\dagger} = J_{\mp}$$

Define simultaneous eigenstates of J^2 and J_z :

$$J^2 |\lambda, m\rangle = \lambda \hbar^2 |\lambda, m\rangle \qquad J_z |\lambda, m\rangle$$

 $\mathbf{n} \cdot \mathbf{J} = n_x J_x + n_y J_y + n_z J_z$

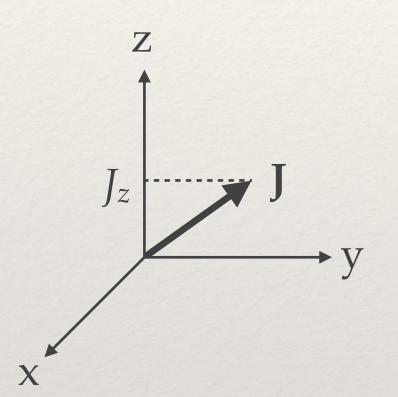
 $[J_7, J^2] = 0$

$$[J_z, J_{\pm}] = \pm \hbar J_{\pm} \qquad [J^2, J_{\pm}] = 0$$

 $J_{\pm} | \lambda, m \rangle \propto | \lambda, m \pm 1 \rangle$ $m\rangle = m\hbar |\lambda, m\rangle$



What is the physical meaning of λ ? Is it just total AM squared? We have:



 $m^2 \leq \lambda$

This follows from: $\langle \lambda, m | J_x^2 + J_y^2$

"z-component squared cannot be more than total AM squared"

 $\langle \lambda, m | J'$

We can guess the physical meaning of *m*: it is essentially the *z*-component of AM.

$$\begin{aligned} \langle \lambda, m \rangle &\geq 0 \iff \langle \lambda, m | J^2 - J_z^2 | \lambda, m \rangle \geq 0 \\ 2^2 - J_z^2 | \lambda, m \rangle &= \langle \lambda, m | \lambda \hbar^2 - m^2 \hbar^2 | \lambda, m \rangle \\ &= (\lambda - m^2) \hbar^2 \langle \lambda, m | \lambda, m \rangle \\ &= (\lambda - m^2) \hbar^2 \langle \lambda, m | \lambda, m \rangle \end{aligned}$$



Therefore there must exist a *maximum* value of *m*: call it *j*. We must have:

$$J_+ |\lambda, j\rangle = 0$$
 Otherwise

Now calculate:

$$J_{-}J_{+} = (J_{x} - iJ_{y})(J_{x} + iJ_{y}) = J_{x}^{2} + J_{y}^{2} + i[J_{x}, J_{y}] = J^{2} - J_{z}^{2} - \hbar J_{z}$$

 $J_{-}J_{+}|\lambda,j\rangle = 0 \implies (J^{2} - J_{z}^{2} - \hbar J_{z}^{2})$

$$\Rightarrow (\lambda - j^2 - j)/$$

se we create a state having $m^2 = (j+1)^2 > \lambda$.

$$J_z)|\lambda,j\rangle = 0$$

 $h^2 |\lambda, j\rangle = 0 \qquad \Rightarrow \lambda = j(j+1)$

Similarly, there must exist a *minimum* value of *m*: call it *j*'. We must have:

$$J_{-}|\lambda,j'
angle=0$$
 Otherwis

Now calculate:

$$J_{+}J_{-} = (J_{x} + iJ_{y})(J_{x} - iJ_{y}) = J_{x}^{2} + J_{y}^{2} - i[J_{x}, J_{y}] = J^{2} - J_{z}^{2} + \hbar J_{z}$$

 $J_{+}J_{-}|\lambda,j'\rangle = 0 \implies (J^{2} - J_{7}^{2} + \hbar J_{7})|\lambda,j'\rangle = 0$

$$\Rightarrow (\lambda - j'^2 + j')$$

se we create a state having $m^2 = (j'-1)^2 > \lambda$.

 $h^{2}|\lambda,j'\rangle = 0 \implies \lambda = j'(j'-1)$

These two values of λ must be self-consistent:

$$\lambda = j(j + 1) = j'(j' - 1)$$

This implies:

$$j' = -j$$
 or

valid

We can therefore exactly compute λ :

The quantity *j* is more fundamental than λ , so we will re-label our states as:

j' = j + 1

violates our inequality, so discard it.

$\lambda = i(i+1)$

We'll use this notation from now on: $\lambda, m \rightarrow j, m \rightarrow j, m \rightarrow Note: J^2 | j, m \rightarrow j (j + 1)\hbar^2 | j, m \rightarrow j$



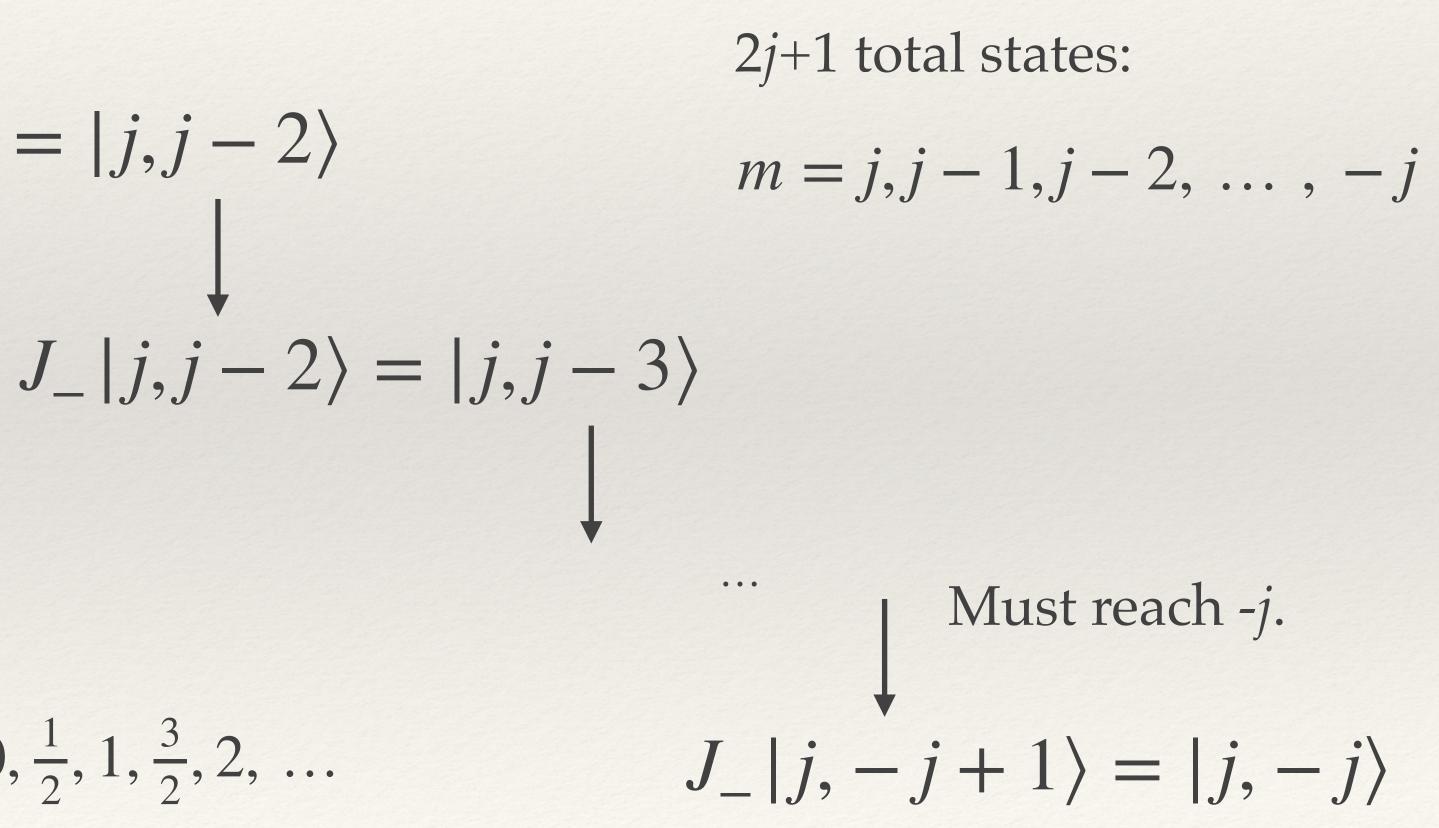
Quantized values

What are the allowed values of *j* and *m*? Start with m = j and work down:

$$J_{-}|j,j\rangle = |j,j-1\rangle$$

$$J_{-}|j,j-1\rangle = |j,j-1\rangle$$

Since we take an integer number of steps, the distance j-(-j) = 2j must be an integer. Allowed values for *j* are: $j = 0, \frac{1}{2}, 1, \frac{3}{2}, 2, ...$



Matrix elements

What do the raising and lowering operators look like in this basis? We know: $J_{+}|j,m\rangle = c_{+}\hbar|j,m+1\rangle$

$$\langle j, m | J_J_+ | j, m \rangle = c_+^* c_+ \hbar^2 \langle j, m+1 | j, m+1 \rangle = |c_+|^2 \hbar^2$$

We also have:

$$\langle j,m | J_J_+ | j,m \rangle = \langle j,m | J^2 - J_z^2 - \hbar J_z | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m^2 - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m | j,m \rangle = \left(j(j+1) - m \right) \hbar^2 \langle j,m \rangle = \left$$

We can always make a choice of phase such that: $c_{+} =$

And similarly for the lowering operator we find:

$$c_{-} = \sqrt{j(j)}$$

and
$$J_{\perp}|j,m\rangle = c_{\hbar}|j,m-1\rangle$$

$$(+1) - m(m+1)$$

(i + 1) - m(m - 1)



Matrix elements

Putting these expressions together we find:

$$\langle j, m' | J_+ | j, m \rangle = c_+ \hbar \langle j, m | j, m + 1 \rangle = \sqrt{j(j+1) - m(m+1)} \hbar \delta_{m', m+1}$$

And similarly for the lowering operator we find:

$$\langle j, m' | J_{-} | j, m \rangle = c_{-} \hbar \langle j, m | j, m - 1 \rangle = \sqrt{j(j+1) - m(m-1)} \hbar \delta_{m',m-1}$$

From these expressions, we can write explicit matrices that act on the space of spin-*j* states.