

*Prof. Steven Flammia*

---

# Quantum Mechanics

---

## Lecture 2

Time evolution and the Schrödinger equation;  
The Hamiltonian as the generator of time translations;  
Wave functions in infinite-dimensional Hilbert spaces;  
Position operator and position basis.





# Commuting operators

---

Consider the case of two **nondegenerate** operators  $A$  and  $B$   
Suppose they are Hermitian and that they commute.

More generally, commuting Hermitian operators share a common eigenbasis.  
(The proof can be done by generalizing the above argument.)

To track commutativity (or lack thereof), introduce the **commutator**:

Many nice algebraic identities...



# Unitary time evolution

---

Let's look at the unitary operator that translates a state in time:

Recall, it must be unitary to conserve probability.

Rather than study the most general such operator, Taylor expand for small time:

Unitarity at first order in  $dt$  implies:



---

# The Schrödinger equation

---

What about at large times? We can expand again, but around  $t$ .



Schrödinger equation,  
operator form

When  $H$  is time-independent, the general solution is:



# The Schrödinger equation

---

Applying both sides to some initial state  $|\psi(0)\rangle$ , we find



Schrödinger equation,  
state vector form

$$i\hbar \frac{d}{dt} U(t) = H U(t)$$

Schrödinger equation,  
operator form

Is this still unitary for all  $t$ , not just  $dt$ ? Assuming  $H$  is independent of  $t$ :



---

# The Hamiltonian operator

---

Let's continue assuming that  $H$  is independent of  $t$ .

Recall that  $H$  has units of energy.

It commutes with  $U(t)$ :

It is self-adjoint, so it is an observable with real eigenvalues.

What is the expected value of  $H$ ?

We therefore define  $H$  to be the Hamiltonian or energy operator.



---

# The Hamiltonian operator

---

What are the eigenstates of  $H$ ? The *energy eigenstates*:

The energy eigenstates are “stationary” with respect to time:

Superpositions of energy eigenstates have non-trivial dynamics.

*Example:*



---

# Time dependence of expected values

---

What about time dependence of expected values more generally?

Operators  $A$  that are independent of time are conserved iff they commute with  $H$ .



# Position basis

---

Our derivation of the Schrödinger equation was completely general. But let's focus on a special case more challenging than spin degrees of freedom: *position*.

Unlike spin, which takes a finite set of values, position is a *continuous* variable.

In analogy with spin, let's consider a 1D line and define a position operator:

We should be able to expand any state in the *position basis*. Because position is a continuous variable, the resolution of the identity takes an integral form:



# Real-space wave functions

---

This suggests defining the *wave function*  $\psi(x)$ :

What is the Born rule probability for finding the particle at  $x$ ?

What is the Born rule probability for finding the particle between  $x$  and  $x+dx$ ?



---

# Position eigenstates and the Dirac delta function

---

Are the eigenstates of the position operator valid wave functions?

The Dirac delta function:

The Dirac delta function is **not** a normalizable wave function, so position eigenstates are not physically realizable.



---

# Expected values and overlaps

---

The trick is always to insert a resolution of the identity.