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# Quantum Mechanics

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## Lecture 4

Particle in a box;  
Wave mechanics in 3D;  
The generator of rotations;  
Angular momentum.





# A quick recap

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Translation is generated by momentum:

$$T(\delta x) = 1 - \frac{i}{\hbar} \hat{p} \delta x \quad T(a) = e^{-i\hat{p}a/\hbar}$$

Momentum obeys the relations:

$$[\hat{x}, \hat{p}] = i\hbar \quad \hat{p} = \hat{p}^\dagger \quad \hat{p} \xrightarrow{\text{x basis}} \frac{\hbar}{i} \frac{\partial}{\partial x}$$

The time-independent Schrödinger equation for a 1D particle with potential is:

$$-\frac{2m(E - V(x))}{\hbar^2} \psi_E(x) = \frac{\partial^2}{\partial x^2} \psi_E(x)$$

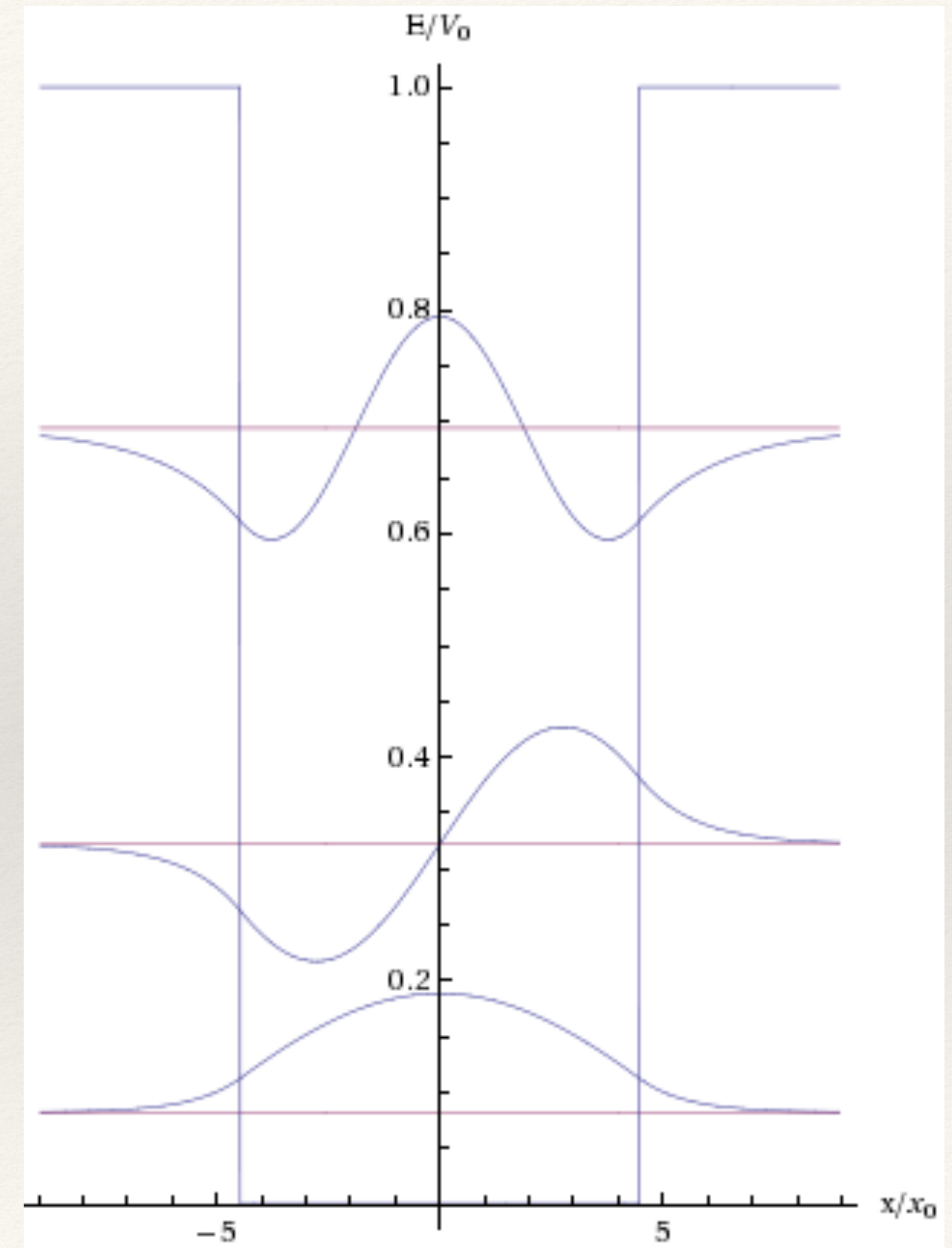


# Particle in a box / finite square well

These tools give us insight into interesting physics!

$$-\frac{2m(E - V(x))}{\hbar^2}\psi_E(x) = \frac{\partial^2}{\partial x^2}\psi_E(x) \quad V(x) = \begin{cases} 0 & |x| < L/2 \\ V_0 & \text{otherwise} \end{cases}$$

- ❖ Fixed number of bound states
- ❖ Quantized energy levels,  $E_n \propto n^2/mL^2$  when  $V_0 \rightarrow \infty$
- ❖ Particles “exist” in classically forbidden regions
- ❖ Energy eigenstates with  $E > V_0$  cannot be normalized
- ❖ Scattering states must form wave packets
- ❖ States can scatter *back* off a potential well





# Towards 3D wave mechanics: rotations

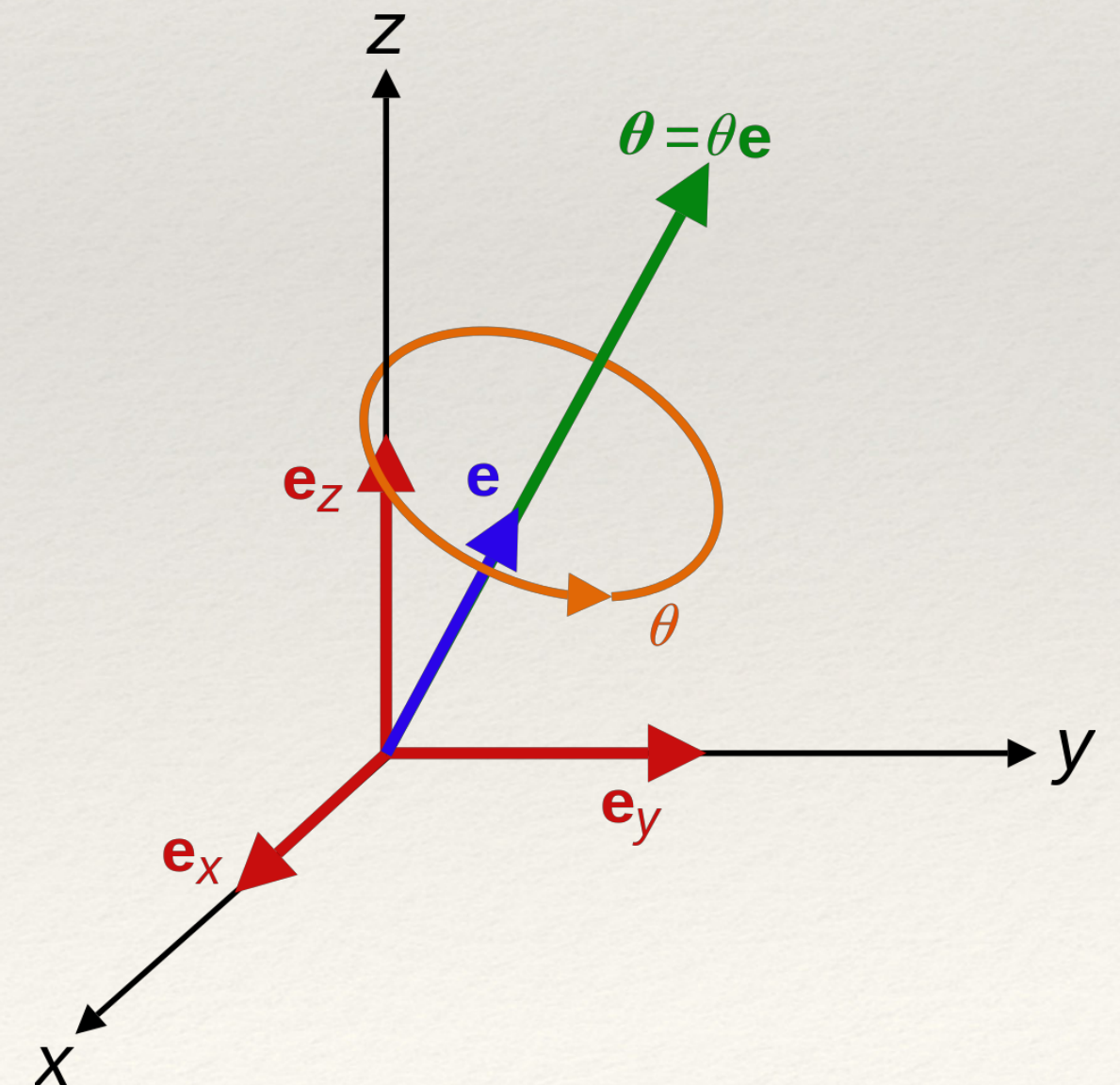
Our eventual goal is to get a quantitative understanding of the physics of atoms. “Particle in a box” is too simple to give predictions that match experiment.

To achieve that, we’ll need to build on our toolkit from 1D and incorporate one more phenomenon that doesn’t exist in only one dimension: **rotations**.

Introduce the rotation operator:

How do rotation operators act on spin states?

*Example:* spin-1/2





# Rotation operators

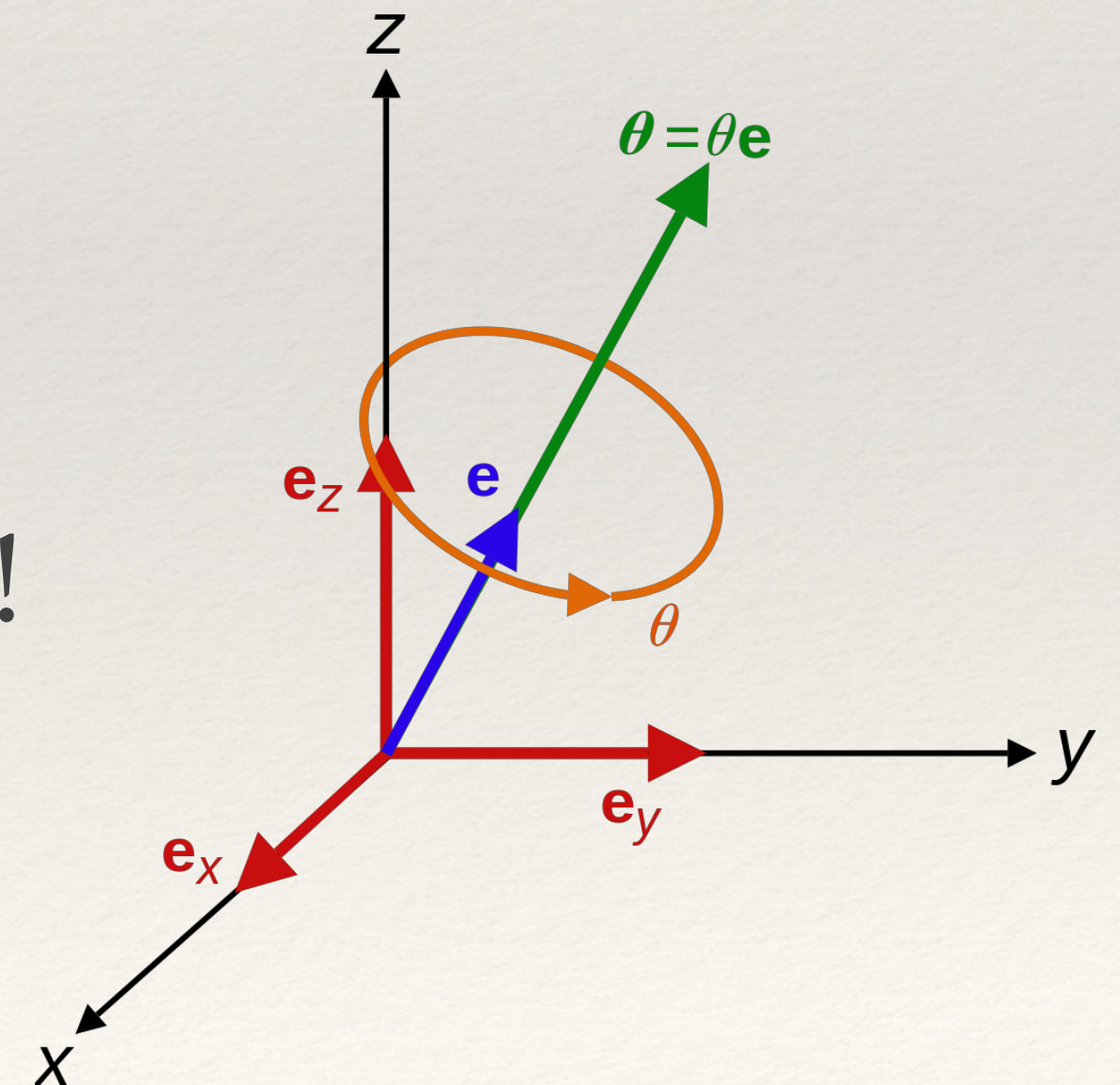
Let's be more systematic and recycle the recipe for time and space translations. Introduce a generator for infinitesimal rotations around each axis:

It must be unitary, so  $J_z$  is Hermitian.

And  $J_z$  clearly has units of angular momentum.

Moreover, z-aligned states must be eigenstates for every  $\theta$ !

*Example:* spin-1/2

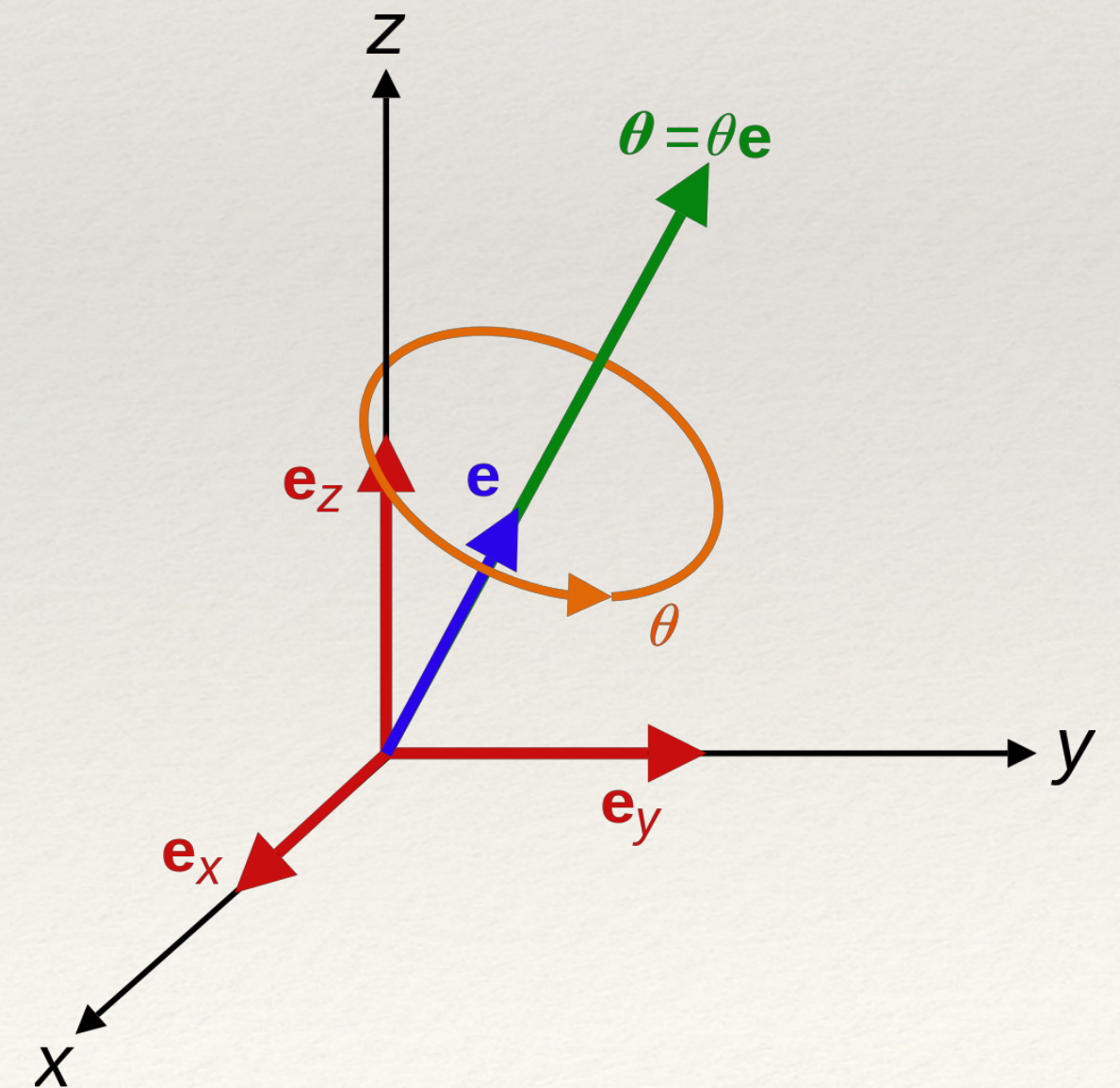




# Rotation operators

How does the phase depend on the angle? To remain consistent with interpretation as angular momentum, we must have for spin-1 / 2:

Sanity check:





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# Commutation relations for rotation operators

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More generally, we can have a rotation about any axis:

All the above insights can be generalized to any spin, not just spin  $1/2$ .  
To go beyond the spin- $1/2$ , let us study the commutation relations



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# Total angular momentum: the Casimir operator

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Rotations about different axes don't commute, but there is another invariant.

Recall:



# Simultaneous eigenstates, raising and lowering operators

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Since  $J^2$  and  $J_z$  are commuting and self-adjoint, they have a common eigenbasis.

To be more explicit, we must define raising and lowering operators:

Do these commute with  $J_z$ ?



# Simultaneous eigenstates, raising and lowering operators

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Although they don't commute, the mutual action on the basis is key:

Similarly:

Note that raising and lowering commute with  $J^2$ , so  $\lambda$  is unchanged.

These operators add and subtract one quanta of angular momentum to the z-projection of the angular momentum eigenstates.